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MINIMAL RIESZ ENERGY ON THE SPHERE FOR AXIS-SUPPORTED EXTERNAL FIELDS

J. S. BRAUCHART, P. D. DRAGNEV†, AND E. B. SAFF‡

ABSTRACT. We investigate the minimal Riesz s -energy problem for positive measures on the d -dimensional unit sphere \mathbb{S}^d in the presence of an external field induced by a point charge, and more generally by a line charge. The model interaction is that of Riesz potentials $|\mathbf{x} - \mathbf{y}|^{-s}$ with $d - 2 \leq s < d$. For a given axis-supported external field, the support and the density of the corresponding extremal measure on \mathbb{S}^d is determined. The special case $s = d - 2$ yields interesting phenomena, which we investigate in detail. A weak* asymptotic analysis is provided as $s \rightarrow (d - 2)^+$.

1. INTRODUCTION AND RESULTS

1.1. **Potential-theoretical preliminaries.** Let $\mathbb{S}^d := \{\mathbf{x} \in \mathbb{R}^{d+1} : |\mathbf{x}| = 1\}$ be the unit sphere in \mathbb{R}^{d+1} , where $|\cdot|$ denotes the Euclidean norm, and let $\sigma = \sigma_d$ be the unit Lebesgue surface measure on \mathbb{S}^d . Recall that, using cylindrical coordinates

$$(1.1) \quad \mathbf{x} = \left(\sqrt{1 - u^2} \bar{\mathbf{x}}, u \right), \quad -1 \leq u \leq 1, \quad \bar{\mathbf{x}} \in \mathbb{S}^{d-1},$$

we can write the decomposition

$$(1.2) \quad d\sigma_d(\mathbf{x}) = \frac{\omega_{d-1}}{\omega_d} (1 - u^2)^{d/2-1} du d\sigma_{d-1}(\bar{\mathbf{x}}).$$

Here ω_d is the surface area of \mathbb{S}^d , and the ratio of these areas can be evaluated as

$$(1.3) \quad \frac{\omega_d}{\omega_{d-1}} = \int_{-1}^1 (1 - u^2)^{d/2-1} du = \frac{\sqrt{\pi} \Gamma(d/2)}{\Gamma((d+1)/2)} = 2^{d-1} \frac{[\Gamma(d/2)]^2}{\Gamma(d)}.$$

Given a compact set $E \subset \mathbb{S}^d$, consider the class $\mathcal{M}(E)$ of unit positive Borel measures supported on E . For $0 < s < d$, the *Riesz s -potential* and *Riesz s -energy* of a measure $\mu \in \mathcal{M}(E)$ are given, respectively, by

$$U_s^\mu(\mathbf{x}) := \int k_s(\mathbf{x}, \mathbf{y}) d\mu(\mathbf{y}), \quad \mathcal{I}_s(\mu) := \int \int k_s(\mathbf{x}, \mathbf{y}) d\mu(\mathbf{x}) d\mu(\mathbf{y}),$$

where $k_s(\mathbf{x}, \mathbf{y}) := |\mathbf{x} - \mathbf{y}|^{-s}$ is the so-called *Riesz kernel* (for $s = 0$ we use the logarithmic kernel $k_0(\mathbf{x}, \mathbf{y}) := \log(1/|\mathbf{x} - \mathbf{y}|)$ instead). The *s -energy of E* is $W_s(E) := \inf\{\mathcal{I}_s(\mu) : \mu \in \mathcal{M}(E)\}$ and if $W_s(E)$ is finite, there is a unique measure $\mu_{E,s}$ achieving this

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minimal energy, which is called the s -extremal measure on E . The s -capacity of E is defined as $\text{cap}_s(E) := 1/W_s(E)$ for $s > 0$. (In the logarithmic case $s = 0$ we define $\text{cap}_0(E) := \exp\{-W_0(E)\}$, cf. (1.49) for $E = \mathbb{S}^d$.) A property is said to hold *quasi-everywhere* (q.e.) if the exceptional set has s -capacity zero. For more details see [14, Chapter II]. We remind the reader that the s -energy of \mathbb{S}^d is given by

$$(1.4) \quad W_s(\mathbb{S}^d) = \frac{\Gamma(d)\Gamma((d-s)/2)}{2^s \Gamma(d/2)\Gamma(d-s/2)}, \quad 0 < s < d.$$

The weighted s -energy associated with a non-negative lower semi-continuous external field $Q : E \rightarrow [0, \infty]$ is given by

$$(1.5) \quad \mathcal{I}_Q(\mu) := \mathcal{I}_s(\mu) + 2 \int Q(\mathbf{x}) \, d\mu(\mathbf{x}).$$

A measure $\mu_Q \in \mathcal{M}(E)$ such that $\mathcal{I}_Q(\mu_Q) = V_Q$, where

$$(1.6) \quad V_Q := \inf \{ \mathcal{I}_Q(\mu) : \mu \in \mathcal{M}(E) \},$$

is called an *extremal (or positive equilibrium) measure on E associated with $Q(\mathbf{x})$* . The measure μ_Q is characterized by the Gauss variational inequalities

$$(1.7) \quad U_s^{\mu_Q}(\mathbf{x}) + Q(\mathbf{x}) \geq F_Q \quad \text{q.e. on } E,$$

$$(1.8) \quad U_s^{\mu_Q}(\mathbf{x}) + Q(\mathbf{x}) \leq F_Q \quad \text{everywhere on } \text{supp}(\mu_Q),$$

where

$$(1.9) \quad F_Q := V_Q - \int Q(\mathbf{x}) \, d\mu_Q(\mathbf{x}).$$

For simplicity, we suppressed in some of the above notation the dependency on s ; that is, $\mathcal{I}_Q = \mathcal{I}_{Q,s}$, $\mu_Q = \mu_{Q,s}$, etc. We note that for suitable external fields (e.g. continuous on $E = \mathbb{S}^d$), the inequality in (1.7) holds everywhere, which implies that equality holds in (1.8).

The existence, uniqueness, and characterization-related questions concerning extremal potentials with external fields in the most general setting can be found in [23]–[25]. We remark that the logarithmic potential with external fields is treated in depth in [21].

When $Q \equiv 0$ and $\text{cap}_s(E) > 0$, the extremal measure μ_Q is the same as the measure $\mu_E = \mu_{E,s}$.

In [4] *Riesz external fields*

$$(1.10) \quad Q_{\mathbf{a},q}(\mathbf{x}) := Q_{\mathbf{a},q,s}(\mathbf{x}) := q |\mathbf{x} - \mathbf{a}|^{-s} \text{ on } E = \mathbb{S}^d, \quad d-2 < s < d,$$

were considered, where $q > 0$ and \mathbf{a} is a fixed point on \mathbb{S}^d .¹ The motivation for that investigation was to obtain new separation results for minimal s -energy points on the sphere. In the current work we extend that investigation to Riesz external fields $Q_{\mathbf{a},q}$ with $\mathbf{a} \notin \mathbb{S}^d$ and develop a technique for finding the extremal measure associated with more general axis-supported external fields.

¹The case $d = 1$, $s = 0$, where \mathbf{a} is a point on the unit circle was investigated in [15].

1.2. Signed Equilibrium. We note that for $d = 2$ and $s = 1$ it is a standard electrostatic problem to find the charge density (signed measure) on a charged, insulated, conducting sphere in the presence of a point charge q placed off the sphere (see [9, Chapter 2]). This motivates us to give the following definition (see [3]).

Definition 1.1. Given a compact subset $E \subset \mathbb{R}^p$ ($p \geq 3$) and an external field Q , we call a signed measure $\eta_{E,Q} = \eta_{E,Q,s}$ supported on E and of total charge $\eta_{E,Q}(E) = 1$ a *signed s -equilibrium on E associated with Q* if its weighted Riesz s -potential is constant on E , that is

$$(1.11) \quad U_s^{\eta_{E,Q}}(\mathbf{x}) + Q(\mathbf{x}) = F_{E,Q} \quad \text{for all } \mathbf{x} \in E.$$

The choice of the normalization $\eta_{E,Q}(E) = 1$ is just for convenience in the applications here. Lemma 2.1 below establishes that if a signed s -equilibrium $\eta_{E,Q}$ exists, then it is unique.

In [6] Fabrikant et al give a derivation of certain signed Riesz equilibria on suitably parametrized surfaces in \mathbb{R}^3 , including spherical caps when $Q(\mathbf{x}) \equiv 0$. We remark that the determination of signed equilibria is a substantially easier problem than that of finding non-negative extremal measures, which is the goal of this paper. However, the solution to the former problem is useful in solving the latter problem.

Our first result establishes existence of the signed s -equilibrium associated with the Riesz external field $Q_{\mathbf{a},q}$, $\mathbf{a} \notin \mathbb{S}^d$, defined in (1.10). We assume that \mathbf{a} lies above the North Pole $\mathbf{p} := (\mathbf{0}, 1)$, that is $\mathbf{a} = (\mathbf{0}, R)$ and $R > 1$ (the case $R < 1$ is handled by inversion).

Throughout, ${}_2F_1\left(\begin{smallmatrix} a, b \\ c \end{smallmatrix}; z\right)$ and ${}_2\tilde{F}_1\left(\begin{smallmatrix} a, b \\ c \end{smallmatrix}; z\right)$ denote the Gauss hypergeometric function and its regularized form ² with series expansions

$$(1.12) \quad {}_2F_1\left(\begin{smallmatrix} a, b \\ c \end{smallmatrix}; z\right) := \sum_{n=0}^{\infty} \frac{(a)_n (b)_n}{(c)_n} \frac{z^n}{n!}, \quad {}_2\tilde{F}_1\left(\begin{smallmatrix} a, b \\ c \end{smallmatrix}; z\right) := \sum_{n=0}^{\infty} \frac{(a)_n (b)_n}{\Gamma(n+c)} \frac{z^n}{n!}, \quad |z| < 1,$$

where $(a)_0 := 1$ and $(a)_n := a(a+1) \cdots (a+n-1)$ for $n \geq 1$ is the Pochhammer symbol. The incomplete Beta function and the Beta function are defined as

$$(1.13) \quad B(x; \alpha, \beta) := \int_0^x v^{\alpha-1} (1-v)^{\beta-1} dv, \quad B(\alpha, \beta) := B(1; \alpha, \beta),$$

whereas the regularized incomplete Beta function is given by

$$(1.14) \quad I(x; a, b) := B(x; a, b) / B(a, b).$$

Theorem 1.2. *Let $0 < s < d$ and $R > 1$. The signed s -equilibrium $\eta_{\mathbf{a}} = \eta_{\mathbb{S}^d, Q_{\mathbf{a},q}, s}$ on \mathbb{S}^d associated with the Riesz external field $Q_{\mathbf{a},q}$, $\mathbf{a} = R\mathbf{p}$, is given by*

$$(1.15) \quad d\eta_{\mathbf{a}}(\mathbf{x}) = \left\{ 1 + \frac{qU_s^\sigma(\mathbf{a})}{W_s(\mathbb{S}^d)} - \frac{q(R^2-1)^{d-s}}{W_s(\mathbb{S}^d)|\mathbf{x}-\mathbf{a}|^{2d-s}} \right\} d\sigma(\mathbf{x}).$$

Furthermore, $U_s^\sigma(\mathbf{a}) = \int k_s(\mathbf{a}, \mathbf{y}) d\sigma(\mathbf{y})$ has the following representation:

$$(1.16) \quad U_s^\sigma(\mathbf{a}) = (R+1)^{-s} {}_2F_1\left(\begin{smallmatrix} s/2, d/2 \\ d \end{smallmatrix}; \frac{4R}{(R+1)^2}\right).$$

²which is well-defined even for c a negative integer

We remark that in the Coulomb case $d = 2$ and $s = 1$, the representation (1.15) is well-known from elementary physics (cf. [9, p. 61]).

The next result explicitly shows the relationship between q and R so that $\mu_{Q_{\mathbf{a},q}}$ coincides with the signed equilibrium and has as support the entire sphere.

Corollary 1.3. *Let $0 < s < d$ and $R = |\mathbf{a}| > 1$. Then $\text{supp}(\mu_{Q_{\mathbf{a},q}}) = \mathbb{S}^d$ if and only if*

$$(1.17) \quad \frac{W_s(\mathbb{S}^d)}{q} \geq \frac{(R+1)^{d-s}}{(R-1)^d} - U_s^\sigma(\mathbf{a})$$

$$(1.18) \quad = \frac{1}{(R+1)^s} \sum_{k=0}^{\infty} \left[1 - \frac{(s/2)_k}{(d)_k} \right] \frac{(d/2)_k}{k!} \left[\frac{4R}{(R+1)^2} \right]^k.$$

In such a case, $\mu_{Q_{\mathbf{a},q}} = \eta_{\mathbf{a}}$.

Remark 1.4. Observe that the function of R in (1.18) is strictly decreasing for $R > 1$. Thus, for any fixed charge q there is a critical R_q given by equality in (1.17), such that for $R \geq R_q$ the extremal support is the entire sphere.

1.3. The Newtonian case $s = d - 1$. The following example deals with the classical case of a Newtonian potential (relative to the manifold dimension). The example answers a question of A. A. Gonchar; namely, how far from the unit sphere should a unit point charge be placed so that the support of the extremal measure associated with the external field exerted by the charge be the entire sphere?

Example 1.5. Let $d \geq 2$, $s = d - 1$, $q = 1$ and $\mathbf{a} = (\mathbf{0}, R)$. Then $W_s(\mathbb{S}^d) = 1$ (cf. (1.4)) and from the mean-value property for harmonic functions we can write

$$U_s^\sigma(\mathbf{a}) = \frac{1}{R^{d-1}} \quad \text{for } R \geq 1.$$

Thus (1.18) in this case is equivalent to the inequality

$$(1.19) \quad 1 \geq \frac{R+1}{(R-1)^d} - \frac{1}{R^{d-1}} \quad \text{or} \quad 1 \geq \frac{\rho+2}{\rho^d} - \frac{1}{(\rho+1)^{d-1}},$$

where ρ measures the distance between the unit charge and the surface of the sphere. Equality holds, if ρ is an algebraic number satisfying

$$(1.20) \quad P(d; \rho) := (\rho^d - 2 - \rho)(\rho + 1)^{d-1} + \rho^d = 0,$$

or on expanding the polynomial $P(d; \rho)$,

$$(1.21) \quad \sum_{m=0}^{d-1} \binom{d-1}{m} \rho^{m+d} - \sum_{m=0}^{d-1} \left[\binom{d}{m} + \binom{d-1}{m} \right] \rho^m = 0.$$

The monic polynomial³ $P(d; \rho)$ with integer coefficients has odd degree $2d - 1$. Furthermore, $P(d; 1) < 0$ and hence $P(d; \rho)$ has at least one positive root; but, by Descartes' Sign Rule, this is the only positive root. This simple root ρ_+ must be in the interval $(1, 2]$, since $P(d; \rho) > 0$ for $\rho > 2$. Asymptotic analysis shows that

$$(1.22) \quad \rho_+ = 1 + (\log 3)/d + \mathcal{O}(1/d^2) \quad \text{as } d \rightarrow \infty.$$

³Properties of these polynomials will be investigated in a future publication.

Of particular interest is the case when $d = 2$. Then one easily computes that the distance between the point charge and the surface of the sphere is given precisely by the golden ratio

$$\rho_+ = (1 + \sqrt{5})/2.$$

We note that the fact that the inequality $R - 1 \geq \rho_+$ implies $\text{supp}(\mu_{Q_{\mathbf{a},1}}) = \mathbb{S}^2$ follows from an elementary physics argument.

1.4. The Mhaskar-Saff \mathcal{F}_s -functional and the extremal support. An important tool in our analysis is the Riesz analog of the *Mhaskar-Saff F -functional* from classical logarithmic potential in the plane (see [16] and [21, Chapter IV, p. 194]).

Definition 1.6. Given a compact subset $K \subset \mathbb{S}^d$ of positive s -capacity, we define the \mathcal{F}_s -functional of the set K as

$$(1.23) \quad \mathcal{F}_s(K) := W_s(K) + \int Q(\mathbf{x}) \, d\mu_K(\mathbf{x}),$$

where $W_s(K)$ is the s -energy of K and μ_K is the s -extremal measure (without external field) on K .

Remark 1.7. We caution the reader that (1.23) is the negative of the F -functional defined in [16] and [21].

Remark 1.8. When $d - 2 < s < d$, there is a remarkable relationship between the signed equilibrium and the \mathcal{F}_s -functional. Namely, if the signed s -equilibrium on a compact set K associated with Q exists, then $\mathcal{F}_s(K) = F_{K,Q}$, where $F_{K,Q}$ is the constant from (1.11). Indeed, if $\eta_{K,Q}$ exists, we integrate (1.11) with respect to μ_K and interchange the order of integration to obtain the asserted equality.

Remark 1.9. With the notion of the functional \mathcal{F}_s at hand we can restate the results of Theorem 1.2 and Corollary 1.3 as follows: For $0 < s < d$ and $R > 1$ the signed s -equilibrium $\eta_{\mathbf{a}} = \eta_{\mathbb{S}^d, Q_{\mathbf{a},q},s}$ on \mathbb{S}^d associated with $Q_{\mathbf{a},q}$, $\mathbf{a} = R\mathbf{p}$, is given by

$$(1.24) \quad d\eta_{\mathbf{a}}(\mathbf{x}) = \frac{1}{W_s(\mathbb{S}^d)} \left\{ \mathcal{F}_s(\mathbb{S}^d) - q(R^2 - 1)^{d-s} / |\mathbf{x} - \mathbf{a}|^{2d-s} \right\} d\sigma(\mathbf{x}).$$

Moreover, $\text{supp}(\mu_{Q,s}) = \mathbb{S}^d$ (that is $\mu_{Q,s} = \eta_{\mathbf{a}}$) if and only if

$$(1.25) \quad \mathcal{F}_s(\mathbb{S}^d) \geq q(R+1)^{d-s} / (R-1)^d.$$

The following optimization property is the main motivation for introducing the \mathcal{F}_s -functional.

Theorem 1.10. *Let $d - 2 \leq s < d$ with $s > 0$ and Q be an external field on \mathbb{S}^d . Then the \mathcal{F}_s -functional is minimized for $S_Q := \text{supp}(\mu_Q)$.*

The next theorem provides sufficient conditions on a general external field Q that guarantee that the extremal support S_Q is a spherical zone or a spherical cap.

Theorem 1.11. *Let $d - 2 \leq s < d$ with $s > 0$ and the external field $Q : \mathbb{S}^d \rightarrow [0, \infty]$ be rotationally invariant about the polar axis; that is, $Q(\mathbf{z}) = f(\xi)$, where ξ is the altitude of $\mathbf{z} = (\sqrt{1 - \xi^2} \bar{\mathbf{z}}, \xi)$ (see (1.1)). Suppose that f is a convex function on $[-1, 1]$. Then the support of the s -extremal measure μ_Q on \mathbb{S}^d is a spherical zone; namely, there are numbers $-1 \leq t_1 \leq t_2 \leq 1$ such that*

$$(1.26) \quad \text{supp}(\mu_Q) = \Sigma_{t_1, t_2} := \{(\sqrt{1 - u^2} \bar{\mathbf{x}}, u) : t_1 \leq u \leq t_2, \bar{\mathbf{x}} \in \mathbb{S}^{d-1}\}.$$

Moreover, if additionally f is increasing, then $t_1 = -1$ and the support of μ_Q is a spherical cap centered at the South Pole.

It is easy to see that the external field $Q_{\mathbf{a},q}(\mathbf{z}) = q|1 - 2R\xi + R^2|^{-s/2}$ is rotationally invariant about the polar axis and is an increasing and convex function of the altitude ξ of \mathbf{z} . Therefore, from Theorem 1.11 we conclude that the support of the extremal measure $\mu_{Q_{\mathbf{a},q}}$ on \mathbb{S}^d is a spherical cap. In view of Theorem 1.10 we thus need only to minimize the \mathcal{F}_s -functional over the collection of spherical caps centered at the South Pole in order to determine S_Q . For this purpose, in consideration of Remark 1.8, we first seek an explicit representation for the signed equilibria for these spherical caps.

Denote by Σ_t the spherical cap centered at the South Pole

$$(1.27) \quad \Sigma_t := \Sigma_{-1,t},$$

(cf. (1.26)), and let η_t be the signed s -equilibrium on Σ_t associated with $Q_{\mathbf{a},q}$. Using M. Riesz's approach to s -balayage as presented in [14, Chapter IV], we introduce the following s -balayage measures onto Σ_t :

$$(1.28) \quad \epsilon_t = \epsilon_{t,s} := \text{Bal}_s(\delta_{\mathbf{a}}, \Sigma_t), \quad \nu_t = \nu_{t,s} := \text{Bal}_s(\sigma, \Sigma_t),$$

where $\delta_{\mathbf{a}}$ is the unit Dirac-delta measure at \mathbf{a} . Recall that given a measure ν and a compact set K (of the sphere \mathbb{S}^d), the balayage measure $\hat{\nu} := \text{Bal}_s(\nu, K)$ preserves the Riesz s -potential of ν onto the set K and diminishes it elsewhere (on the sphere \mathbb{S}^d). We remark that in what follows an important role is played by the function

$$(1.29) \quad \Phi_s(t) := W_s(\mathbb{S}^d) (1 + q \|\epsilon_t\|) / \|\nu_t\|, \quad d - 2 < s < d.$$

The next assertion is an immediate consequence of the definition of the balayage measures in (1.28). In Lemmas 3.1 and 3.2 below we present explicit formulas for their densities. Their norms are calculated in Lemmas 4.2 and 4.1, respectively. Below we combine these formulas to give an explicit form for the density of the signed s -equilibrium. The only statements requiring further proof is the formula for the weighted s -potential (1.32) when $\xi > t$. We shall do this in Section 6.

Theorem 1.12. *Let $d - 2 < s < d$. The signed s -equilibrium η_t on the spherical cap $\Sigma_t \subset \mathbb{S}^d$ associated with $Q_{\mathbf{a},q}$ is given by*

$$(1.30) \quad \eta_t = \frac{1 + q \|\epsilon_t\|}{\|\nu_t\|} \nu_t - q \epsilon_t.$$

It is absolutely continuous in the sense that for $\mathbf{x} = (\sqrt{1 - u^2} \bar{\mathbf{x}}, u) \in \Sigma_t$,

$$d \eta_t(\mathbf{x}) = \eta'_t(u) \frac{\omega_{d-1}}{\omega_d} (1 - u^2)^{d/2-1} du d \sigma_{d-1}(\bar{\mathbf{x}}),$$

where (with $R = |\mathbf{a}|$ and $r = \sqrt{R^2 - 2Rt + 1}$)

$$\begin{aligned} \eta'_t(u) = & \frac{1}{W_s(\mathbb{S}^d)} \frac{\Gamma(d/2)}{\Gamma(d-s/2)} \left(\frac{1-t}{1-u} \right)^{d/2} \left(\frac{t-u}{1-t} \right)^{(s-d)/2} \\ & \times \left\{ \Phi_s(t) {}_2\tilde{F}_1 \left(\begin{matrix} 1, d/2 \\ 1 - (d-s)/2 \end{matrix}; \frac{t-u}{1-u} \right) \right. \\ & \left. - \frac{q(R+1)^{d-s}}{r^d} {}_2\tilde{F}_1 \left(\begin{matrix} 1, d/2 \\ 1 - (d-s)/2 \end{matrix}; \frac{(R-1)^2}{r^2} \frac{t-u}{1-u} \right) \right\}. \end{aligned}$$

Furthermore, if $\mathbf{z} = (\sqrt{1 - \xi^2} \bar{\mathbf{z}}, \xi) \in \mathbb{S}^d$, the weighted s -potential is given by

$$(1.31) \quad U_s^{\eta_t}(\mathbf{z}) + Q_{\mathbf{a},s}(\mathbf{z}) = \Phi_s(t), \quad \mathbf{z} \in \Sigma_t,$$

$$(1.32) \quad \begin{aligned} U_s^{\eta_t}(\mathbf{z}) + Q_{\mathbf{a},s}(\mathbf{z}) &= \Phi_s(t) + q \frac{1}{\rho^s} \mathrm{I} \left(\frac{(R+1)^2 \xi - t}{r^2} \frac{d-s}{1+\xi}; \frac{d-s}{2}, \frac{s}{2} \right) \\ &\quad - \Phi_s(t) \mathrm{I} \left(\frac{\xi - t}{1+\xi}; \frac{d-s}{2}, \frac{s}{2} \right), \quad \mathbf{z} \in \mathbb{S}^d \setminus \Sigma_t, \end{aligned}$$

where $\rho = \sqrt{R^2 - 2R\xi + 1}$ and $\mathrm{I}(x; a, b)$ is the regularized incomplete Beta function (see (1.14)).

The corresponding statement for the case $s = d - 2$ is given in Theorem 1.16.

Remark 1.13. According to Remark 1.8 we have from Theorem 1.12 that $\mathcal{F}_s(\Sigma_t) = \Phi_s(t)$. Concerning the minimization of this function, we derive the following result.

Theorem 1.14. *Let $d - 2 < s < d$. For the external field $Q_{\mathbf{a},q}(\mathbf{x})$, $\mathbf{a} = (\mathbf{0}, R)$, $R > 1$, the function $\Phi_s(t)$ has precisely one global minimum $t_0 \in (-1, 1]$. This minimum is either the unique solution $t_0 \in (-1, 1)$ of the equation*

$$(1.33) \quad \Phi_s(t) = \frac{q(R+1)^{d-s}}{(R^2 - 2Rt + 1)^{d/2}},$$

or $t_0 = 1$ when such a solution does not exist. Moreover, $t_0 = \max\{t : \eta_t \geq 0\}$.

The extremal measure $\mu_{Q_{\mathbf{a},q}}$ on \mathbb{S}^d is given by η_{t_0} (see (1.30)), and has as support the spherical cap Σ_{t_0} .

Note that, in view of formulas (4.1) and (4.3) for $\|\epsilon_t\|$ and $\|\nu_t\|$ given below, equation (1.33) can be written in terms of hypergeometric functions.

Remark 1.15. The restriction on the parameter s arises in the process of applying the balayage method and the principle of domination. It is a topic for further investigation to extend the range of s for which the conclusion of Theorem 1.14 remains true.

Figure 1 gives an overview of the qualitative behavior of the weighted s -potential of the signed s -equilibrium measure η_t on \mathbb{S}^d associated with the external field Q and its density with respect to $\sigma_d|_{\Sigma_t}$ for s in the range $d - 2 < s < d$ and the choices $t < t_0$, $t = t_0$ and $t > t_0$. We remark that the derivative with respect to ξ of the weighted s -potential becomes $\pm\infty$ as $\xi \rightarrow t^+$ for $t \neq t_0$ and vanishes for $t = t_0 < 1$ (cf. Remark 6.1).

1.5. The exceptional case $s = d - 2$. In this case M. Riesz's approach [14, Chapter IV] has to be modified. Somewhat surprisingly it turns out, as shown in Lemmas 7.2 and 7.5, that the s -balayage measures from (1.28)

$$(1.34) \quad \bar{\epsilon}_t := \epsilon_{t,d-2} = \mathrm{Bal}_{d-2}(\delta_{\mathbf{a}}, \Sigma_t), \quad \bar{\nu}_t := \nu_{t,d-2} = \mathrm{Bal}_{d-2}(\sigma, \Sigma_t)$$

exist and both have a component that is uniformly distributed on the boundary of Σ_t . Moreover, unlike the case $d - 2 < s < d$, the density for $\mu_{Q_{\mathbf{a},q}}$, where $s = d - 2$, does not vanish on the boundary of its support. More precisely, on setting

$$\beta_t(\mathbf{x}) := \delta_t(u) \cdot \sigma_{d-1}(\bar{\mathbf{x}}), \quad \mathbf{x} = (\sqrt{1 - u^2} \bar{\mathbf{x}}, u),$$

we obtain the following result.

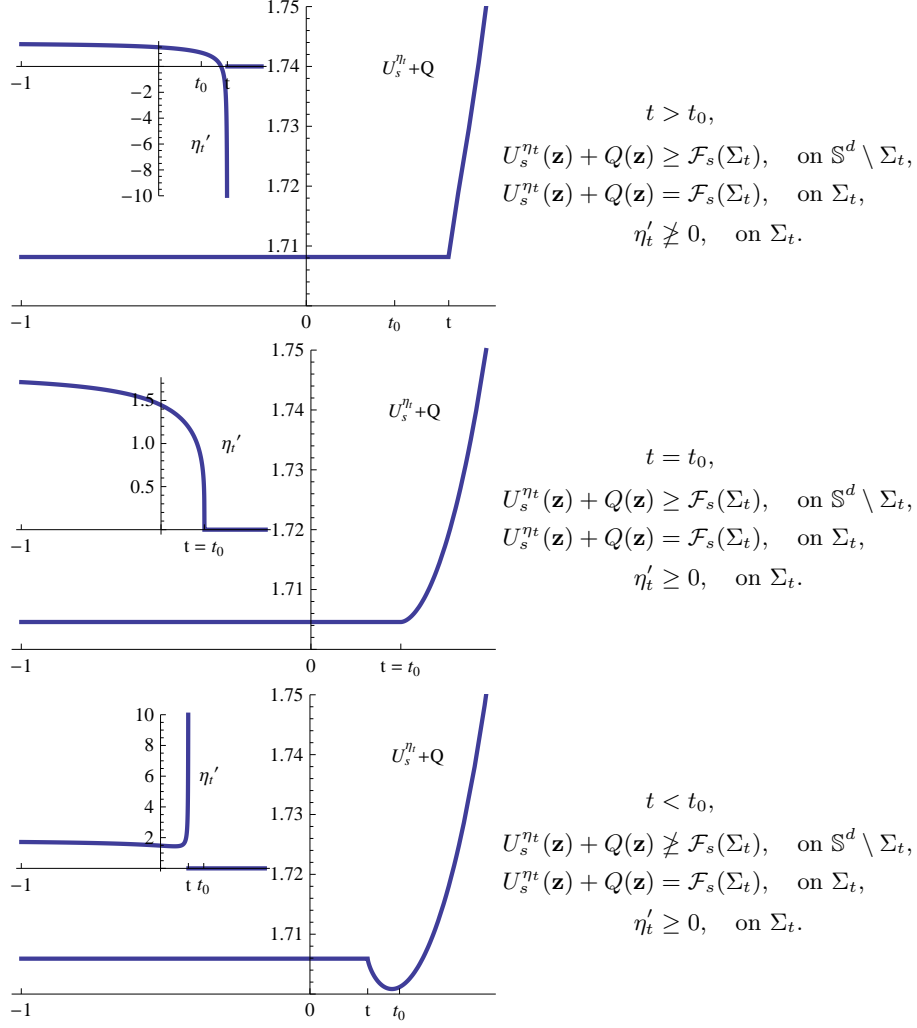


FIGURE 1. The weighted s -potential of η_t for different choices of t ($t > t_0$, $t = t_0$, and $t < t_0$) versus altitude ξ of \mathbf{z} for $d = 2$, $s = 1/2$, $q = 1$, and $R = 3/2$, cf. Theorems 1.12 and 1.14. Insets show the respective density η'_t .

Theorem 1.16. *Let $d \geq 3$. The signed s -equilibrium $\bar{\eta}_t$ on the spherical cap Σ_t associated with $\bar{Q}_{\mathbf{a},q}(\mathbf{x}) = q|\mathbf{x} - \mathbf{a}|^{2-d}$ is given by*

$$\bar{\eta}_t = \frac{\bar{\Phi}_{d-2}(t)}{W_s(\mathbb{S}^d)} \bar{\nu}_t - q\bar{\epsilon}_t, \quad \bar{\Phi}_{d-2}(t) := W_s(\mathbb{S}^d) \frac{1 + q \|\bar{\epsilon}_t\|}{\|\bar{\nu}_t\|},$$

where $\bar{\nu}_t$ and $\bar{\epsilon}_t$ are given in (1.34) and can be written as

$$(1.35) \quad \begin{aligned} d\bar{\eta}_t(\mathbf{x}) = & \frac{1}{W_{d-2}(\mathbb{S}^d)} \left[\bar{\Phi}_{d-2}(t) - \frac{q(R^2-1)^2}{(R^2-2Ru+1)^{d/2+1}} \right] d\sigma_d|_{\Sigma_t}(\mathbf{x}) \\ & + \frac{1-t}{2} (1-t^2)^{d/2-1} \left[\bar{\Phi}_{d-2}(t) - \frac{q(R+1)^2}{(R^2-2Rt+1)^{d/2}} \right] d\beta_t(\mathbf{x}). \end{aligned}$$

For any fixed $t \in (-1, 1)$, the following weak* convergence holds:

$$(1.36) \quad \nu_{t,s} \xrightarrow{*} \bar{\nu}_t, \quad \epsilon_{t,s} \xrightarrow{*} \bar{\epsilon}_t, \quad \text{as } s \rightarrow (d-2)^+.$$

The function $\bar{\Phi}_{d-2}(t)$ has precisely one global minimum $t_0 \in (-1, 1]$. This minimum is either the unique solution $t_0 \in (-1, 1)$ of the equation

$$(1.37) \quad \bar{\Phi}_{d-2}(t) = \frac{q(R+1)^2}{(R^2-2Rt+1)^{d/2}},$$

or $t_0 = 1$ when such a solution does not exist. Moreover, $t_0 = \max\{t : \bar{\eta}_t \geq 0\}$.

The extremal measure $\mu_{\bar{Q}_{\mathbf{a},q}}$ on \mathbb{S}^d is given by

$$(1.38) \quad d\mu_{\bar{Q}_{\mathbf{a},q}}(\mathbf{x}) = d\bar{\eta}_{t_0}(\mathbf{x}) = \frac{\bar{\Phi}_{d-2}(t_0)}{W_{d-2}(\mathbb{S}^d)} \left[1 - \frac{(R-1)^2 (R^2-2Rt_0+1)^{d/2}}{(R^2-2Ru+1)^{d/2+1}} \right] d\sigma_d|_{\Sigma_{t_0}}(\mathbf{x}),$$

and has as support the spherical cap Σ_{t_0} .

In Lemmas 7.2 and 7.5 we give the s -potentials of the balayage measures $\bar{\nu}_t$ and $\bar{\epsilon}_t$ from which the weighted s -potential of $\bar{\eta}_t$ at every $\mathbf{z} \in \mathbb{S}^d$ can be easily obtained.

Remark 1.17. As can be seen from (1.35), depending on the sign of the coefficient of β_t , the signed s -equilibrium $\bar{\eta}_t$ has positive or negative charge on $\partial\Sigma_t$ unless $t = t_0$, in which case the charge on the boundary disappears (see Figure 2).

Next, we describe the results when $d = 2$ and $s = 0$. The external field in this case is $\bar{Q}(\mathbf{x}) = \bar{Q}_{\mathbf{a},q}(\mathbf{x}) = q \log(1/|\mathbf{x} - \mathbf{a}|)$. The total mass of the balayage measures $\bar{\nu}_{t,0}$ and $\bar{\epsilon}_{t,0}$ is preserved, so $\|\bar{\nu}_{t,0}\| = \|\bar{\epsilon}_{t,0}\| = 1$. Thus, the function $\bar{\Phi}_{d-2}(t)$ reduces to $\bar{\Phi}_{d-2}(t) = 1 + q$. The Mhaskar-Saff functional $\mathcal{F}_0(K)$ from (1.23), now defined for compact sets $K \subset \mathbb{S}^2$ with positive logarithmic capacity $\text{cap}_0(K) = \exp\{-W_0(K)\}$, uses the logarithmic energy

$$W_0(K) = \lim_{s \rightarrow 0^+} \frac{dW_s(K)}{ds} \Big|_{s=0}.$$

However, $\mathcal{F}_0(\Sigma_t)$ is no longer equal to $\bar{\Phi}_{d-2}(t)$ (cf. Remark 1.13 and Lemma 7.9). For $K = \mathbb{S}^2$ we have $W_0(\mathbb{S}^2) = 1/2 - \log 2 < 0$. Since Theorem 1.11 can be extended to $s = 0$ if $d = 2$, we deduce that $S_{\bar{Q}} := \text{supp}(\mu_{\bar{Q}})$ will be a spherical cap Σ_{t_0} . Direct calculations show that the Mhaskar-Saff functional \mathcal{F}_0 for spherical caps is still minimized for $S_{\bar{Q}}$. Figure 2 shows the qualitative behavior for the weighted potential in the logarithmic case. (Note, that for $t \neq t_0$ the tangent line to the graph of the weighted logarithmic potential at $\xi \rightarrow t^+$ is **not** vertical like in the case $d - 2 < s < d$ (cf. Figure 1), but it becomes horizontal if $t = t_0 < 1$.)

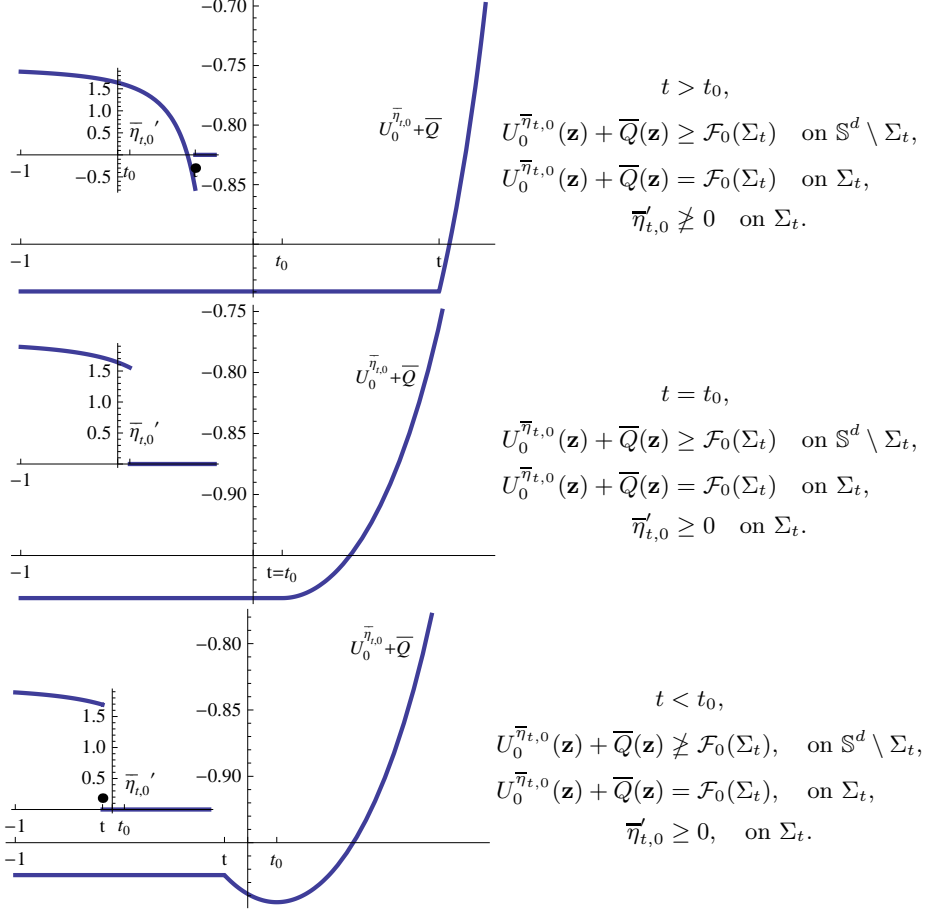


FIGURE 2. The weighted logarithmic potential of $\bar{\eta}_{t,0}$ for different choices of t ($t > t_0$, $t = t_0$, and $t < t_0$) versus altitude ξ of \mathbf{z} for $d = 2$, $s = 0$, $q = 1$, and $R = 2$, cf. Theorem 1.18. Insets show the respective density $\bar{\eta}'_{t,0}$. The black dot indicates the component on the boundary.

Theorem 1.18. *Let $d = 2$ and $s = 0$. The signed s -equilibrium $\bar{\eta}_{t,0}$ on the spherical cap Σ_t associated with $\bar{Q}_{\mathbf{a},q}(\mathbf{x}) = q \log(1/|\mathbf{x} - \mathbf{a}|)$ is given by*

$$\bar{\eta}_{t,0} = (1 + q) \bar{\nu}_{t,0} - q \bar{\epsilon}_{t,0},$$

where $\bar{\nu}_{t,0} = \text{Bal}_0(\sigma_2, \Sigma_t)$ and $\bar{\epsilon}_{t,0} = \text{Bal}_0(\delta_{\mathbf{a}}, \Sigma_t)$. It can be written as

$$(1.39) \quad \begin{aligned} d\bar{\eta}_{t,0}(\mathbf{x}) = & \left[1 + q - \frac{q(R^2 - 1)^2}{(R^2 - 2Ru + 1)^2} \right] d\sigma_2|_{\Sigma_t}(\mathbf{x}) \\ & + \frac{1-t}{2} \left[1 + q - \frac{q(R+1)^2}{R^2 - 2Rt + 1} \right] d\beta_t(\mathbf{x}), \end{aligned}$$

where $\mathbf{x} = (\sqrt{1 - u^2} \bar{\mathbf{x}}, u)$, $\bar{\mathbf{x}} \in \mathbb{S}^1$.

The weighted logarithmic potential of $\bar{\eta}_{t,0}$ satisfies

$$\begin{aligned} U_0^{\bar{\eta}_{t,0}}(\mathbf{z}) + \bar{Q}_{\mathbf{a},q}(\mathbf{z}) &= \mathcal{F}_0(\Sigma_t), & \mathbf{z} \in \Sigma_t, \\ U_0^{\bar{\eta}_{t,0}}(\mathbf{z}) + \bar{Q}_{\mathbf{a},q}(\mathbf{z}) &= \mathcal{F}_0(\Sigma_t) + \frac{1}{2} \log \frac{1+t}{1+\xi} + \frac{q}{2} \log \frac{R^2 - 2Rt + 1}{R^2 - 2R\xi + 1}, & \mathbf{z} \in \mathbb{S}^2 \setminus \Sigma_t, \end{aligned}$$

where $\mathcal{F}_0(\Sigma_t)$ is given below in Lemma 7.9.

The Mhaskar-Saff functional \mathcal{F}_0 is minimized for Σ_{t_0} , where either $t_0 \in (-1, 1]$ is the unique solution of the equation

$$(1.40) \quad 1 + q = q(R+1)^2 / (R^2 - 2Rt + 1),$$

or $t_0 = 1$ if such a solution does not exist. Moreover, $t_0 = \max\{t : \bar{\eta}_{t,0} \geq 0\}$.

The logarithmic extremal measure $\mu_{\bar{Q}_{\mathbf{a},q}}$ on \mathbb{S}^2 is given by

$$(1.41) \quad d\mu_{\bar{Q}_{\mathbf{a},q}}(\mathbf{x}) = d\bar{\eta}_{t_0,0}(\mathbf{x}) = \left[1 + q - \frac{q(R^2 - 1)^2}{(R^2 - 2Ru + 1)^2} \right] d\sigma_2|_{\Sigma_{t_0}}(\mathbf{x}),$$

and has as support the spherical cap Σ_{t_0} .

Remark 1.19. Given R and q , relation (1.40) immediately enables us to find the support Σ_{t_0} of the logarithmic extremal measure $\mu_{\bar{Q}_{\mathbf{a},q}}$ on \mathbb{S}^2 :

$$(1.42) \quad t_0 = \min \left\{ 1, \frac{(R+1)^2}{2R(1+q)} - 1 \right\} = \min \left\{ 1, 1 - \frac{4qR - (R-1)^2}{2R(1+q)} \right\}.$$

Remark 1.20. In general, the density $\bar{\eta}'_{t_0,0}(u)$ in (1.41) does not vanish on the boundary of Σ_{t_0} . In fact, if $t_0 \in (-1, 1)$, then

$$(1.43) \quad \lim_{u \rightarrow t_0} \bar{\eta}'_{t_0,0}(u) = 1 + q - \frac{q(R^2 - 1)^2}{(R^2 - 2Rt_0 + 1)^2} = \frac{1+q}{q} \cdot \frac{4qR - (R-1)^2}{(R+1)^2} > 0.$$

1.6. Axis-supported external fields. It is well known that the balayage of a measure can be represented as a superposition of balayages of Dirac-delta measures. Using this, we extend our results to external fields that are axis-supported s -potentials.

Definition 1.21. We call an external field Q *positive-axis supported*, if

$$(1.44) \quad Q(\mathbf{x}) = \int \frac{d\lambda(R)}{|\mathbf{x} - R\mathbf{p}|^s}, \quad \mathbf{x} \in \mathbb{S}^d,$$

for some finite positive measure λ supported on a compact subset of $(0, \infty)$.

Remark 1.22. Since

$$(1.45) \quad \text{Bal}_s(\delta_{(1/R)\mathbf{p}}, \mathbb{S}^d) = R^s \text{Bal}_s(\delta_{R\mathbf{p}}, \mathbb{S}^d),$$

we can restrict ourselves to measures λ with support in $[1, \infty)$. It is possible to generalize the setting to fields supported on both the negative and positive polar axis as well. This generalization shall be reserved for a later occasion.

We begin with a result that establishes the existence of the signed equilibrium measure $\tilde{\eta}_\lambda$ on \mathbb{S}^d associated with the axis-supported external field Q . Furthermore, a necessary and sufficient condition for coincidence of $\tilde{\eta}_\lambda$ and the extremal measure μ_Q on \mathbb{S}^d is given.

Theorem 1.23. *Let $0 < s < d$ and Q be as in (1.44) with $\text{supp}(\lambda) \subset [1, \infty)$. Then*

$$(1.46) \quad d\tilde{\eta}_\lambda(\mathbf{x}) = \frac{1}{W_s(\mathbb{S}^d)} \left\{ \mathcal{F}_s(\mathbb{S}^d) - \int \frac{(R^2 - 1)^{d-s}}{(R^2 - 2Ru + 1)^{d-s/2}} d\lambda(R) \right\} d\sigma(\mathbf{x}).$$

Moreover, $\text{supp}(\mu_Q) = \mathbb{S}^d$ (that is $\mu_Q = \tilde{\eta}_\lambda$) if and only if

$$(1.47) \quad \mathcal{F}_s(\mathbb{S}^d) \geq \int \frac{(R+1)^{d-s}}{(R-1)^d} d\lambda(R).$$

The above result, appropriately modified, also holds for the logarithmic case. We shall use the Mhaskar-Saff functional for the logarithmic case

$$(1.48) \quad \mathcal{F}_0(K) := W_0(K) + \int Q(\mathbf{x}) d\mu_{K,0}(\mathbf{x}),$$

where K is a compact subset of \mathbb{S}^d with finite logarithmic energy $W_0(K)$ and $\mu_{K,0}$ is the logarithmic extremal measure on K (without external field). In particular,

$$(1.49) \quad W_0(\mathbb{S}^d) = \lim_{s \rightarrow 0^+} \frac{dW_s(\mathbb{S}^d)}{ds} \Big|_{s=0} = -\log 2 - \frac{1}{2} \psi(d/2) - \frac{1}{2} \psi(d),$$

where $\psi(z) := \Gamma'(z)/\Gamma(z)$ denotes the digamma function.

Theorem 1.24. *Let $d = 2$, $s = 0$, and Q be the positive-axis supported external field*

$$(1.50) \quad Q(\mathbf{z}) = \int \log \frac{1}{|\mathbf{z} - \mathbf{a}|} d\lambda(R)$$

with $\text{supp}(\lambda) \subset [1, \infty)$. Then the signed logarithmic equilibrium measure $\tilde{\eta}_{\lambda,0}$ on \mathbb{S}^d associated with Q is given by

$$(1.51) \quad d\tilde{\eta}_{\lambda,0}(\mathbf{x}) = \left\{ 1 + \|\lambda\| - \int \frac{(R^2 - 1)^d}{(R^2 - 2Ru + 1)^d} d\lambda(R) \right\} d\sigma(\mathbf{x}).$$

Its weighted logarithmic potential is given by

$$(1.52) \quad U_0^{\tilde{\eta}_{\lambda,0}}(\mathbf{z}) + Q(\mathbf{z}) = \mathcal{F}_0(\mathbb{S}^d), \quad \mathbf{z} \in \mathbb{S}^d.$$

Moreover, $\text{supp}(\mu_Q) = \mathbb{S}^d$ (that is the logarithmic extremal measure μ_Q on \mathbb{S}^d coincides with $\tilde{\eta}_{\lambda,0}$) if and only if

$$(1.53) \quad 1 + \|\lambda\| \geq \int \frac{(R+1)^d}{(R-1)^d} d\lambda(R).$$

The next assertion deals with the signed equilibrium measure $\tilde{\eta}_t$ on a spherical cap Σ_t for Q of the form (1.44).

Theorem 1.25. *Let $d - 2 < s < d$ and Q be as in (1.44) with $\text{supp}(\lambda) \subset [1, \infty)$. The signed s -equilibrium $\tilde{\eta}_t$ on the spherical cap Σ_t associated with Q is given by*

$$(1.54) \quad \tilde{\eta}_t = \frac{\tilde{\Phi}_s(t)}{W_s(\mathbb{S}^d)} \nu_t - \tilde{\epsilon}_t, \quad \tilde{\Phi}_s(t) := W_s(\mathbb{S}^d) \frac{1 + \|\tilde{\epsilon}_t\|}{\|\nu_t\|},$$

where ν_t is defined in (1.28) and

$$(1.55) \quad \tilde{\epsilon}_t := \text{Bal}_s(\lambda, \Sigma_t) = \int \text{Bal}_s(\delta_{R\mathbf{p}}, \Sigma_t) d\lambda(R).$$

The signed s -equilibrium $\tilde{\eta}_t$ can be written as

$$d\tilde{\eta}_t(\mathbf{x}) = \tilde{\eta}'_t(u, R) \frac{\omega_{d-1}}{\omega_d} (1-u^2)^{d/2-1} du d\sigma_{d-1}(\bar{\mathbf{x}}), \quad \mathbf{x} \in \Sigma_t,$$

where

$$\begin{aligned} \tilde{\eta}'_t(u, R) &= \frac{1}{W_s(\mathbb{S}^d)} \frac{\Gamma(d/2)}{\Gamma(d-s/2)} \left(\frac{1-t}{1-u} \right)^{d/2} \left(\frac{t-u}{1-t} \right)^{(s-d)/2} \\ &\times \left\{ \tilde{\Phi}_s(t) {}_2\tilde{F}_1 \left(\begin{matrix} 1, d/2 \\ 1 - (d-s)/2 \end{matrix}; \frac{t-u}{1-u} \right) \right. \\ &\quad \left. - \int \frac{(R+1)^{d-s}}{(R^2-2Rt+1)^{d/2}} {}_2\tilde{F}_1 \left(\begin{matrix} 1, d/2 \\ 1 - (d-s)/2 \end{matrix}; \frac{(R-1)^2}{R^2-2Rt+1} \frac{t-u}{1-u} \right) d\lambda(R) \right\}. \end{aligned}$$

Furthermore, the function $\tilde{\Phi}_s$ has precisely one global minimum in $(-1, 1]$. This minimum is either the unique solution $t_\lambda \in (-1, 1)$ of the equation

$$(1.56) \quad \tilde{\Phi}_s(t) = \int \frac{(R+1)^{d-s}}{(R^2-2Rt+1)^{d/2}} d\lambda(R),$$

or $t_\lambda = 1$ when such a solution does not exist. Moreover, $t_\lambda := \max\{t : \tilde{\eta}_t \geq 0\}$, $\mu_Q = \tilde{\eta}_{t_\lambda}$, and $\text{supp}(\mu_Q) = \Sigma_{t_\lambda}$, where μ_Q is the extremal measure on \mathbb{S}^d associated with Q .

Theorem 1.25 can be also extended to the case $s = d-2$ and $d \geq 3$. We present

Theorem 1.26. *Let $s = d-2$, $d \geq 3$ and Q be as in (1.44) with $\text{supp}(\lambda) \subset [1, \infty)$. The signed s -equilibrium $\tilde{\eta}_t$ on the spherical cap Σ_t associated with Q is given by*

$$(1.57) \quad \tilde{\eta}_t = \frac{\tilde{\Phi}_{d-2}(t)}{W_s(\mathbb{S}^d)} \bar{\nu}_t - \tilde{\epsilon}_t, \quad \tilde{\Phi}_{d-2}(t) := W_s(\mathbb{S}^d) \frac{1 + \|\tilde{\epsilon}_t\|}{\|\bar{\nu}_t\|},$$

where $\bar{\nu}_t$ is defined in (1.34) and

$$(1.58) \quad \tilde{\epsilon}_t := \text{Bal}_{d-2}(\lambda, \Sigma_t) = \int \text{Bal}_{d-2}(\delta_{R\mathbf{p}}, \Sigma_t) d\lambda(R).$$

The signed s -measure $\tilde{\eta}_t$ can be written as

$$(1.59) \quad d\tilde{\eta}_t(\mathbf{x}) = g(u) d\sigma_d|_{\Sigma_t}(\mathbf{x}) + h(u) d\beta_t(\mathbf{x}),$$

where, when using Lemmas 7.2 and 7.5, we have for $-1 \leq u \leq t$

$$(1.60) \quad g(u) = \frac{1}{W_{d-2}(\mathbb{S}^d)} \left[\tilde{\Phi}_{d-2}(t) - \int \frac{(R^2-1)^2}{(R^2-2Ru+1)^{d/2+1}} d\lambda(R) \right],$$

$$(1.61) \quad h(u) = \frac{1-t}{2} \left[\tilde{\Phi}_{d-2}(t) - \int \frac{(R+1)^2}{(R^2-2Rt+1)^{d/2}} d\lambda(R) \right] (1-t^2)^{d/2-1}.$$

For any fixed $t \in (-1, 1)$, the following weak* convergence holds:

$$(1.62) \quad \tilde{\epsilon}_{t,s} \xrightarrow{*} \tilde{\epsilon}_t \quad \text{as } s \rightarrow (d-2)^+.$$

The function $\tilde{\Phi}_{d-2}$ has precisely one global minimum in $(-1, 1]$. This minimum is either the unique solution $t_\lambda \in (-1, 1)$ of the equation

$$(1.63) \quad \tilde{\Phi}_{d-2}(t) = \int \frac{(R+1)^2}{(R^2-2Rt+1)^{d/2}} d\lambda(R),$$

or $t_\lambda = 1$ when such a solution does not exist. Moreover, $t_\lambda := \max\{t : \tilde{\eta}_t \geq 0\}$, $\mu_Q = \tilde{\eta}_{t_\lambda}$, and $\text{supp}(\mu_Q) = \Sigma_{t_\lambda}$, where μ_Q is the extremal measure on \mathbb{S}^d .

Next, we describe the results when $d = 2$ and $s = 0$. The external field in this case is

$$(1.64) \quad \tilde{Q}(\mathbf{x}) = \tilde{Q}_{\mathbf{a},q}(\mathbf{x}) = \int \log \frac{1}{|\mathbf{x} - R\mathbf{p}|} d\lambda(R), \quad \mathbf{x} \in \mathbb{S}^2,$$

for some finite positive measure λ supported on a compact subset of $[1, \infty)$.

We show a result, which generalizes Theorem 1.18.

Theorem 1.27. *Let $d = 2$ and $s = 0$. Let \tilde{Q} be as in (1.64) with $\text{supp}(\lambda) \subset [1, \infty)$. The signed logarithmic equilibrium $\tilde{\eta}_{t,0}$ on the spherical cap Σ_t associated with \tilde{Q} is given by*

$$(1.65) \quad \tilde{\eta}_{t,0} = (1 + \|\lambda\|) \bar{\nu}_{t,0} - \tilde{\epsilon}_{t,0},$$

where $\bar{\nu}_{t,0} = \text{Bal}_0(\sigma_2, \Sigma_t)$, $\bar{\epsilon}_{t,0} = \text{Bal}_0(\delta_{R\mathbf{p}}, \Sigma_t)$, and

$$(1.66) \quad \tilde{\epsilon}_{t,0} := \text{Bal}_0(\lambda, \Sigma_t) = \int \text{Bal}_0(\delta_{R\mathbf{p}}, \Sigma_t) d\lambda(R).$$

It can be written as

$$(1.67) \quad \begin{aligned} d\tilde{\eta}_{t,0}(\mathbf{x}) &= \left[1 + \|\lambda\| - \int \frac{(R^2 - 1)^2}{(R^2 - 2Ru + 1)^2} d\lambda(R) \right] d\sigma_2|_{\Sigma_t}(\mathbf{x}) \\ &\quad + \frac{1-t}{2} \left[1 + \|\lambda\| - \int \frac{(R+1)^2}{R^2 - 2Rt + 1} d\lambda(R) \right] d\beta_t(\mathbf{x}). \end{aligned}$$

The weighted logarithmic potential of $\tilde{\eta}_{t,0}$ satisfies

$$\begin{aligned} U_0^{\tilde{\eta}_{t,0}}(\mathbf{z}) + \tilde{Q}(\mathbf{z}) &= W_0(\Sigma_t) + \int \tilde{Q} d\mu_{\Sigma_t,0} =: \tilde{\mathcal{F}}_0(\Sigma_t), \quad \mathbf{z} \in \Sigma_t, \\ U_0^{\tilde{\eta}_{t,0}}(\mathbf{z}) + \tilde{Q}(\mathbf{z}) &= \tilde{\mathcal{F}}_0(\Sigma_t) + \frac{1}{2} \log \frac{1+t}{1+\xi} + \int \frac{1}{2} \log \frac{R^2 - 2Rt + 1}{R^2 - 2R\xi + 1} d\lambda(R), \quad \mathbf{z} \in \mathbb{S}^2 \setminus \Sigma_t. \end{aligned}$$

The Mhaskar-Saff functional $\tilde{\mathcal{F}}_0$ (explicitly given in (8.8)) is minimized for Σ_{t_λ} , where either $t_\lambda \in (-1, 1)$ is the unique solution of the equation

$$(1.68) \quad 1 + \|\lambda\| = \int \frac{(R+1)^2}{R^2 - 2Rt + 1} d\lambda(R),$$

or $t_\lambda = 1$ if such a solution does not exist. Moreover, $t_\lambda = \max\{t : \tilde{\eta}_t \geq 0\}$, $\text{supp}(\mu_{\tilde{Q}}) = \Sigma_{t_\lambda}$, and $\mu_{\tilde{Q}} = \tilde{\eta}_{t_\lambda,0}$.

Remark 1.28. In general, the density $\tilde{\eta}'_{t_\lambda,0}(u)$ with respect to $\sigma_2|_{\Sigma_{t_\lambda}}$ of the extremal measure μ_Q on \mathbb{S}^2 in Theorem 1.27 does not vanish on the boundary of Σ_{t_λ} . In fact, if $t_\lambda \in (-1, 1)$, then

$$(1.69) \quad \lim_{u \rightarrow t_\lambda} \tilde{\eta}'_{t_\lambda,0}(u) = 1 + \|\lambda\| - \int \frac{(R^2 - 1)^2 d\lambda(R)}{(R^2 - 2Rt_\lambda + 1)^2} = \int \frac{2R(1 - t_\lambda) d\lambda(R)}{(R^2 - 2Rt_\lambda + 1)^2} > 0.$$

The remainder of this paper is structured as follows. In Section 2 we show the uniqueness of the signed equilibrium and prove Theorem 1.2 and Corollary 1.3. In Section 3 a suitable Kelvin transform of points and measures is considered and explicit formulas for the densities of the measures in (1.30) are found in Lemmas 3.1 and 3.2. The norms of these measures are computed in Section 4. The proofs of Theorems 1.10, 1.11, and 1.14 are given in Section 5. The weighted s -potential of the signed s -equilibrium is given in Section 6. Section 7 considers the special case $s = d - 2$ and the proofs of Theorems 1.16 and 1.18 are provided. Finally, in Section 8 we prove the generalization of the results to axis-supported external fields.

2. SIGNED EQUILIBRIUM ASSOCIATED WITH AN EXTERNAL FIELD

First, we consider some preliminaries on the Kelvin transformation (spherical inversion) of points and measures. Inversion in a sphere is a basic technique in electrostatics (method of electrical images, cf. Jackson [9]) and in general in potential theory (cf. Kellog [10] and Landkof [14]). Kelvin transformation (of a function) is linear, preserves harmonicity (in the classical case), and preserves positivity. We shall make use of this method and of balayage to conveniently infer representations of the signed equilibrium associated with an external field from known results.

2.1. The Kelvin transformation. Let us denote by K_R the Kelvin transformation (stereographic projection) with center $\mathbf{a} = (\mathbf{0}, R)$ and radius $\sqrt{R^2 - 1}$, that is for any point $\mathbf{x} \in \mathbb{R}^{d+1}$ the image $\mathbf{x}^* := K_R(\mathbf{x})$ lies on a ray stemming from \mathbf{a} , and passing through \mathbf{x} such that

$$(2.1) \quad |\mathbf{x} - \mathbf{a}| \cdot |\mathbf{x}^* - \mathbf{a}| = R^2 - 1.$$

Thus, the transformation of the distance is given by the formula

$$(2.2) \quad |\mathbf{x}^* - \mathbf{y}^*| = (R^2 - 1) \frac{|\mathbf{x} - \mathbf{y}|}{|\mathbf{x} - \mathbf{a}| |\mathbf{y} - \mathbf{a}|}, \quad \mathbf{x}, \mathbf{y} \in \mathbb{S}^d.$$

It is easy to see that $K_R(\mathbb{S}^d) = \mathbb{S}^d$, where K_R sends the spherical cap $A_R := \{(\sqrt{1 - u^2} \bar{\mathbf{x}}, u) : 1/R \leq u \leq 1, \bar{\mathbf{x}} \in \mathbb{S}^{d-1}\}$ to $B_R := \{(\sqrt{1 - u^2} \bar{\mathbf{x}}, u) : -1 \leq u \leq 1/R, \bar{\mathbf{x}} \in \mathbb{S}^{d-1}\}$ and vice versa, with the points on the boundary being fixed. In particular, the North Pole $\mathbf{p} = (\mathbf{0}, 1)$ goes to the South Pole $\mathbf{q} := (\mathbf{0}, -1)$. The image of $\mathbf{x} = (\sqrt{1 - u^2} \bar{\mathbf{x}}, u)$ is $\mathbf{x}^* = (\sqrt{1 - (u^*)^2} \bar{\mathbf{x}}, u^*)$, where the relation between u and u^* is given by

$$(2.3) \quad 1 + u^* = \frac{(R + 1)^2}{R^2 - 2Ru + 1} (1 - u).$$

The last equation is derived from the similar triangles proportion

$$|\mathbf{x}^* - \mathbf{q}| / |\mathbf{q} - \mathbf{a}| = |\mathbf{x} - \mathbf{p}| / |\mathbf{x} - \mathbf{a}|$$

and the formulas $|\mathbf{x}^* - \mathbf{q}|^2 = 2(1 + u^*)$, $|\mathbf{x} - \mathbf{p}|^2 = 2(1 - u)$, $|\mathbf{q} - \mathbf{a}| = R + 1$, and $|\mathbf{x} - \mathbf{a}|^2 = R^2 - 2Ru + 1$. Finally, we point out that

$$(2.4) \quad |\mathbf{x}^* - \mathbf{a}|^{-d} d\sigma(\mathbf{x}^*) = |\mathbf{x} - \mathbf{a}|^{-d} d\sigma(\mathbf{x}),$$

which can be easily seen from the relation $(\mathbf{x}^* - \mathbf{a}) / |\mathbf{x}^* - \mathbf{a}| = (\mathbf{x} - \mathbf{a}) / |\mathbf{x} - \mathbf{a}|$.

Next, we recall that given a measure λ with no point mass at \mathbf{a} , its Kelvin transformation (associated with a fixed s) $\lambda^* = \mathcal{K}_{R,s}(\lambda)$ is a measure defined by

$$(2.5) \quad d\lambda^*(\mathbf{x}^*) := \frac{(R^2 - 1)^{s/2}}{|\mathbf{x} - \mathbf{a}|^s} d\lambda(\mathbf{x}).$$

The s -potentials of the two measures are related as follows (see, for example, [4, Section 5, Equation (5.1)])

$$(2.6) \quad U_s^{\lambda^*}(\mathbf{x}^*) = \int \frac{d\lambda^*(\mathbf{y}^*)}{|\mathbf{x}^* - \mathbf{y}^*|^s} = \int \frac{|\mathbf{x} - \mathbf{a}|^s d\lambda(\mathbf{y})}{(R^2 - 1)^{s/2} |\mathbf{x} - \mathbf{y}|^s} = \frac{|\mathbf{x} - \mathbf{a}|^s}{(R^2 - 1)^{s/2}} U_s^\lambda(\mathbf{x}).$$

Note that the Kelvin transformation has the duality property $\mathcal{K}_{R,s}(\lambda^*(\mathbf{x}^*)) = \lambda(\mathbf{x})$.

2.2. Signed equilibrium. We first establish the uniqueness of the signed equilibrium, provided it exists.

Lemma 2.1. *Let $0 \leq s < d$. If a signed s -equilibrium $\eta_{E,Q}$ exists, then it is unique.*

Proof. The lemma easily follows from the positivity of the s -energy of signed measures. Indeed, suppose η_1 and η_2 are two signed s -equilibria on E associated with the same external field Q . Then

$$U_s^{\eta_1}(\mathbf{x}) + Q(\mathbf{x}) = F_1, \quad U_s^{\eta_2}(\mathbf{x}) + Q(\mathbf{x}) = F_2 \quad \text{for all } \mathbf{x} \in E.$$

Subtracting the two equations and integrating with respect to $\eta_1 - \eta_2$ we obtain

$$\mathcal{I}_s(\eta_1 - \eta_2) = \int [U_s^{\eta_1}(\mathbf{x}) - U_s^{\eta_2}(\mathbf{x})] d(\eta_1 - \eta_2)(\mathbf{x}) = 0,$$

and from [14, Theorem 1.15] we conclude that $\eta_1 = \eta_2$ (see also [7, Section 5]). When $d = 2$ and $s = 0$ instead of [14, Theorem 1.15] we could use [22, Theorem 4.1] to prove the assertion of the Lemma. When $d > 2$ and $s = 0$ we could use [19, p. 6]. Note that $\eta_1 - \eta_2$ is the difference of two signed measures with total charge 1. \square

We are now in a position to find the signed equilibrium for the external field $Q_{\mathbf{a},q}$ defined by a point charge q at \mathbf{a} (see (1.10)).

Proof of Theorem 1.2. Let

$$\epsilon_{\mathbf{a}} := \frac{(R^2 - 1)^{d-s}}{W_s(\mathbb{S}^d) |\mathbf{x} - \mathbf{a}|^{2d-s}} d\sigma(\mathbf{x}), \quad \sigma = \sigma_d.$$

We apply the Kelvin transformation (2.1) to the s -potential

$$U_s^{\epsilon_{\mathbf{a}}}(\mathbf{z}) = \int_{\mathbb{S}^d} \frac{(R^2 - 1)^{d-s}}{W_s(\mathbb{S}^d) |\mathbf{z} - \mathbf{x}|^s |\mathbf{x} - \mathbf{a}|^{2d-s}} d\sigma(\mathbf{x}).$$

From (2.2) and (2.4) (recall that $\mathcal{K}_R(\mathbb{S}^d) = \mathbb{S}^d$) we obtain

$$U_s^{\epsilon_{\mathbf{a}}}(\mathbf{z}) = |\mathbf{z} - \mathbf{a}|^{-s} \int_{\mathbb{S}^d} \frac{1}{W_s(\mathbb{S}^d) |\mathbf{z}^* - \mathbf{x}^*|^s} d\sigma(\mathbf{x}^*) = \frac{1}{|\mathbf{z} - \mathbf{a}|^s},$$

where we used that $U_s^\sigma(\mathbf{z}^*) = W_s(\mathbb{S}^d)$ for all $\mathbf{z}^* \in \mathbb{S}^d$. Hence, $\epsilon_{\mathbf{a}} = \epsilon_1$ (see (1.28)). For $\eta_{\mathbf{a}}$ defined in (1.15), we therefore derive

$$U_s^{\eta_{\mathbf{a}}}(\mathbf{z}) + Q_{\mathbf{a},q}(\mathbf{z}) = W_s(\mathbb{S}^d) + qU_s^\sigma(\mathbf{a}), \quad \text{for all } \mathbf{z} \in \mathbb{S}^d.$$

In addition, one similarly finds

$$\int_{\mathbb{S}^d} \frac{(R^2 - 1)^{d-s}}{|\mathbf{x} - \mathbf{a}|^{2d-s}} d\sigma(\mathbf{x}) = \int_{\mathbb{S}^d} \frac{1}{|\mathbf{x}^* - \mathbf{a}|^s} d\sigma(\mathbf{x}^*) = U_s^\sigma(\mathbf{a}),$$

and consequently $\eta_{\mathbf{a}}(\mathbb{S}^d) = 1$. Therefore, $\eta_{\mathbf{a}}$ is the required signed s -equilibrium.

Finally, to derive (1.16), using (1.2) and (1.3), we evaluate

$$\begin{aligned} U_s^\sigma(\mathbf{a}) &= \int_{\mathbb{S}^d} \frac{1}{|\mathbf{x} - \mathbf{a}|^s} d\sigma_d(x) = \frac{\omega_{d-1}}{\omega_d} \int_{-1}^1 \frac{(1-u^2)^{d/2-1}}{(R^2 - 2Ru + 1)^{s/2}} du \\ (2.7) \quad &= (R+1)^{-s} {}_2F_1\left(\frac{s}{2}, \frac{d}{2}; \frac{4R}{(R+1)^2}\right). \end{aligned}$$

In the last step we used the standard substitution $2v = 1 + u$ and the integral representation of the hypergeometric function [1, Eq. 15.3.1]. \square

The proof of Corollary 1.3 is an easy consequence of the uniqueness of the extremal measure associated with an external field.

Proof of Corollary 1.3. We observe that the (strictly decreasing) density in (1.15) is at minimum on \mathbb{S}^d when $\mathbf{x} = \mathbf{p}$. So, non-negativity at the North Pole implies that the signed equilibrium is positive everywhere else on \mathbb{S}^d , in which case it coincides with the extremal measure on \mathbb{S}^d . On the other hand, if $\text{supp}(\mu_{Q_{\mathbf{a},q}}) = \mathbb{S}^d$, then the variational inequalities (1.7) and (1.8) yield $\mu_{Q_{\mathbf{a},q}} = \eta_{\mathbf{a}}$; and the density in (1.15) is again non-negative at \mathbf{p} . What remains to show is that (1.17) is equivalent to

$$1 + \frac{qU_s^\sigma(\mathbf{a})}{W_s(\mathbb{S}^d)} - \frac{q(R^2 - 1)^{d-s}}{W_s(\mathbb{S}^d)|\mathbf{p} - \mathbf{a}|^{2d-s}} \geq 0,$$

which can be easily seen by using $|\mathbf{p} - \mathbf{a}| = R - 1$. Finally, using the series expansion of (1.16) and

$$\frac{(R+1)^d}{(R-1)^d} = \left[1 - \frac{4R}{(R+1)^2}\right]^{-d/2} = \sum_{k=0}^{\infty} \frac{(d/2)_k}{k!} \left[\frac{4R}{(R+1)^2}\right]^k,$$

we derive (1.18). \square

3. THE s -BALAYAGE MEASURES ν_t AND ϵ_t

In this section we show that for s in the range $d-2 < s < d$, the measures ν_t and ϵ_t are absolutely continuous with respect to the normalized area surface measure σ_d (restricted to the spherical cap Σ_t) and we find their densities.

3.1. The balayage measures. We now focus on the two balayage measures in (1.28). The second one, ν_t , has already been found in [4, Section 3, Equations (3.19) and (4.6)]. It is an absolutely continuous measure on Σ_t (see (1.27)), given by the following formula:

$$(3.1) \quad d\nu_t(\mathbf{x}) = (1 + J_t(\mathbf{x})) \frac{\omega_{d-1}}{\omega_d} (1 - u^2)^{d/2-1} du d\sigma_{d-1}(\bar{\mathbf{x}}),$$

where

$$J_t(\mathbf{x}) := \frac{1}{\Gamma((d-s)/2)\Gamma(1-(d-s)/2)} \left(\frac{1-t}{1-u}\right)^{d/2} \left(\frac{t-u}{1-t}\right)^{(s-d)/2} \\ \times \int_0^1 v^{d/2-1} (1-v)^{1+(d-s)/2-1} \left(1 - \frac{1-t}{1-u}v\right)^{-1} dv.$$

It is convenient to obtain a closed form for $J_t(\mathbf{x})$ in terms of hypergeometric functions. By [1, Eq. 15.3.1]

$$J_t(\mathbf{x}) := \frac{\Gamma(d/2)\Gamma(1+(d-s)/2)}{\Gamma((d-s)/2)\Gamma(1-(d-s)/2)\Gamma(1+d-s/2)} \\ \times \left(\frac{1-t}{1-u}\right)^{d/2} \left(\frac{t-u}{1-t}\right)^{(s-d)/2} {}_2F_1\left(1, d/2; 1+d-s/2; \frac{1-t}{1-u}\right).$$

The application of [1, Eq. 15.3.6] yields an expansion near $u = t$,

$$J_t(\mathbf{x}) = -1 + \frac{\Gamma(d/2)}{\Gamma(d-s/2)} \left(\frac{1-t}{1-u}\right)^{d/2} \left(\frac{t-u}{1-t}\right)^{(s-d)/2} {}_2\tilde{F}_1\left(1, d/2; 1-(d-s)/2; \frac{t-u}{1-u}\right).$$

Substituting the last relation into (3.1) and simplifying we derive the following lemma.

Lemma 3.1. *Let $d-2 < s < d$. The measure $\nu_t = \text{Bal}_s(\sigma, \Sigma_t)$ is given by*

$$(3.2) \quad d\nu_t(\mathbf{x}) = \nu'_t(u) \frac{\omega_{d-1}}{\omega_d} (1-u^2)^{d/2-1} du d\sigma_{d-1}(\bar{\mathbf{x}}), \quad \mathbf{x} \in \Sigma_t,$$

where the density $\nu'_t(u)$ is given by

$$(3.3) \quad \nu'_t(u) := \frac{\Gamma(d/2)}{\Gamma(d-s/2)} \left(\frac{1-t}{1-u}\right)^{d/2} \left(\frac{t-u}{1-t}\right)^{(s-d)/2} {}_2\tilde{F}_1\left(1, d/2; 1-(d-s)/2; \frac{t-u}{1-u}\right).$$

To determine the s -balayage ϵ_t , we recall the formulas for the Kelvin transformation of measures and the relation of the corresponding potentials (see (2.5) and (2.6)). Let λ^* be the extremal measure on the set $\Sigma_t^* := \mathcal{K}_R(\Sigma_t)$, normalized so that its potential $U_s^{\lambda^*}(\mathbf{x}^*) = 1$ for $\mathbf{x}^* \in \Sigma_t^*$. Then, using (2.1) and (2.6) we derive just as in [4, Section 3, Equation (3.7)] that

$$(3.4) \quad \epsilon_t(\mathbf{x}) = (R^2 - 1)^{-s/2} \mathcal{K}_{R,s}(\lambda^*(\mathbf{x}^*)).$$

Since the image Σ_t^* of Σ_t is also a spherical cap, this time centered at the North Pole, we can utilize a formula similar to (3.2) for its extremal measure. If $\Sigma_t = \{\mathbf{x} : -1 \leq u \leq t\}$, then $\Sigma_t^* = \{\mathbf{x} : 1 \geq u^* \geq t^*\}$, where u^* and t^* are related to u and t by (2.3). If we set $\nu_t^* := \text{Bal}(\sigma, \Sigma_t^*)$, then $\lambda^* = \nu_t^*/W_s(\mathbb{S}^d)$; hence we get

$$(3.5) \quad d\lambda^*(\mathbf{x}^*) = (\lambda^*)'(u^*) \frac{\omega_{d-1}}{\omega_d} [1 - (u^*)^2]^{d/2-1} du^* d\sigma_{d-1}(\bar{\mathbf{x}}^*),$$

where the density is given by

$$(\lambda^*)'(u^*) := \frac{1}{W_s(\mathbb{S}^d)} \frac{\Gamma(d/2)}{\Gamma(d-s/2)} \left(\frac{1+t^*}{1+u^*}\right)^{d/2} \left(\frac{u^*-t^*}{1+t^*}\right)^{(s-d)/2} \\ \times {}_2\tilde{F}_1\left(1, d/2; 1-(d-s)/2; \frac{u^*-t^*}{1+u^*}\right).$$

(We remark that the last formula (up to a multiplicative constant) for the special case $d = 2$ was first derived by Fabrikant et al [6].) From (2.3) we get

$$(3.6) \quad \frac{1 + u^*}{1 + t^*} = \frac{R^2 - 2Rt + 1}{R^2 - 2Ru + 1} \cdot \frac{1 - u}{1 - t},$$

from which it follows that

$$(3.7) \quad [1 - (u^*)^2]^{d/2-1} du^* = \left(\frac{R^2 - 1}{R^2 - 2Ru + 1} \right)^d (1 - u^2)^{d/2-1} du.$$

Substituting (3.6) and (3.7) in (3.5) and using (3.4) and (2.5) we obtain the next lemma.

Lemma 3.2. *Let $d - 2 < s < d$. The measure $\epsilon_t = \text{Bal}_s(\delta_{\mathbf{a}}, \Sigma_t)$ is given by*

$$(3.8) \quad d\epsilon_t(\mathbf{x}) = \epsilon'_t(u) \frac{\omega_{d-1}}{\omega_d} (1 - u^2)^{d/2-1} du d\sigma_{d-1}(\bar{\mathbf{x}}), \quad \mathbf{x} \in \Sigma_t,$$

and setting $r^2 := R^2 - 2Rt + 1$, the density is given by

$$(3.9) \quad \begin{aligned} \epsilon'_t(u) := & \frac{1}{W_s(\mathbb{S}^d)} \frac{\Gamma(d/2)}{\Gamma(d-s/2)} \frac{(R+1)^{d-s}}{r^d} \left(\frac{1-t}{1-u} \right)^{d/2} \\ & \times \left(\frac{t-u}{1-t} \right)^{(s-d)/2} {}_2\tilde{F}_1 \left(\begin{matrix} 1, d/2 \\ 1 - (d-s)/2 \end{matrix}; \frac{(R-1)^2}{r^2} \frac{t-u}{1-u} \right). \end{aligned}$$

3.2. Positivity of the signed equilibrium of a spherical cap. The following lemma establishes a condition for positivity of the signed equilibrium

$$d\eta_t(\mathbf{x}) = \eta'_t(u) \frac{\omega_{d-1}}{\omega_d} (1 - u^2)^{d/2-1} du d\sigma_{d-1}(\bar{\mathbf{x}}).$$

Lemma 3.3. *Let $d - 2 < s < d$. If for some $\gamma > 0$ we have $\eta'_t(u) \geq 0$ for $u \in (t - \gamma, t)$, then*

$$(3.10) \quad \Phi_s(t) \geq q(R+1)^{d-s}/r^d, \quad r^2 = R^2 - 2Rt + 1,$$

and, consequently, $\eta'_t(u) > 0$ for all $-1 \leq u < t < 1$.

Proof. By equation (1.30) in Theorem 1.12 (which easily follows from the balayage properties), the definition of $\Phi_s(t)$ (cf. (1.29)), and the Lemmas 3.1 and 3.2 we get

$$(3.11) \quad \begin{aligned} \eta'_t(u) = & \frac{1}{W_s(\mathbb{S}^d)} \frac{\Gamma(d/2)}{\Gamma(d-s/2)} \left(\frac{1-t}{1-u} \right)^{d/2} \left(\frac{t-u}{1-t} \right)^{(s-d)/2} \\ & \times \left\{ \Phi_s(t) {}_2\tilde{F}_1 \left(\begin{matrix} 1, d/2 \\ 1 - (d-s)/2 \end{matrix}; \frac{t-u}{1-u} \right) \right. \\ & \left. - \frac{q(R+1)^{d-s}}{r^d} {}_2\tilde{F}_1 \left(\begin{matrix} 1, d/2 \\ 1 - (d-s)/2 \end{matrix}; \frac{(R-1)^2}{r^2} \frac{t-u}{1-u} \right) \right\}. \end{aligned}$$

Using (3.11) and the non-negativity hypothesis for $\eta'_t(u)$, we get

$$\begin{aligned} \lim_{u \rightarrow t^-} \left[(t-u)^{(d-s)/2} \eta'_t(u) \right] &= \frac{1}{W_s(\mathbb{S}^d)} \frac{\Gamma(d/2)}{\Gamma(d-s/2) \Gamma(1 - (d-s)/2)} \\ & \times (1-u)^{(d-s)/2} \left\{ \Phi_s(t) - \frac{q(R+1)^{d-s}}{r^d} \right\} \geq 0. \end{aligned}$$

In particular, the expression in braces is non-negative for $d - 2 < s < d$.

For $R \neq 1$ we have $(R - 1)^2 < r^2$. Thus, the first hypergeometric function in (3.11) is strictly larger than the second one for all $-1 \leq u < t$ and $d - 2 < s < d$. Hence, using $\Phi_s(t) \geq q(R + 1)^{d-s}/r^d$, we have

$$\begin{aligned} \eta'_t(u) &> \frac{1}{W_s(\mathbb{S}^d)} \frac{\Gamma(d/2)}{\Gamma(d-s/2)} \left(\frac{1-t}{1-u} \right)^{d/2} \left(\frac{t-u}{1-t} \right)^{(s-d)/2} \\ &\quad \times {}_2\tilde{F}_1 \left(\begin{matrix} 1, d/2 \\ 1 - (d-s)/2 \end{matrix}; \frac{t-u}{1-u} \right) \left\{ \Phi_s(t) - \frac{q(R+1)^{d-s}}{r^d} \right\} \geq 0, \end{aligned}$$

which shows that $\eta'_t(u) > 0$ for all $-1 \leq u < t$. \square

Remark 3.4. We note that in the limit $R \rightarrow 1$ relation (3.10) becomes the same as in [4, Eq. (5.9)]. It also follows from the proof of Lemma 3.3 that the sign of the difference $\Phi_s(t) - q(R + 1)^{d-s}/r^d$ is determined by the sign of $\eta'_t(u)$ near the boundary of the spherical cap Σ_t , that is for u near t^- , and vice versa.

Remark 3.5. Equality in relation (3.10) yields $\lim_{u \rightarrow t^-} \eta'_t(u) = 0$. This follows from (3.11) and the identity

$$\begin{aligned} &{}_2\tilde{F}_1 \left(\begin{matrix} 1, d/2 \\ 1 - (d-s)/2 \end{matrix}; \frac{t-u}{1-u} \right) - {}_2\tilde{F}_1 \left(\begin{matrix} 1, d/2 \\ 1 - (d-s)/2 \end{matrix}; \frac{(R-1)^2}{r^2} \frac{t-u}{1-u} \right) \\ &= \sum_{n=1}^{\infty} \frac{(d/2)_n}{\Gamma(n+1 - (d-s)/2)} \left\{ 1 - [(R-1)/r]^{2n} \right\} \left(\frac{t-u}{1-u} \right)^n. \end{aligned}$$

4. THE NORMS $\|\epsilon_t\|$ AND $\|\nu_t\|$.

In this section we compute the norms of the measures in (1.30).

Lemma 4.1. *Let $d - 2 < s < d$. Then*

$$(4.1) \quad \|\epsilon_t\| = \frac{2^{1-d} \Gamma(d)}{\Gamma(d-s/2) \Gamma(s/2)} \frac{(R+1)^{d-s}}{W_s(\mathbb{S}^d)} \int_{-1}^t \frac{(1+u)^{s/2-1} (1-u)^{d-s/2-1}}{(R^2 - 2Ru + 1)^{d/2}} du.$$

Proof. From (3.8) and (3.9)

$$\begin{aligned} \|\epsilon_t\| &= \frac{\omega_{d-1}}{\omega_d} \int_{-1}^t \epsilon'_t(u) (1-u^2)^{d/2-1} du = \frac{\Gamma(d/2)}{\Gamma(d-s/2)} \frac{\omega_{d-1}}{\omega_d} \frac{(R+1)^{d-s}}{W_s(\mathbb{S}^d) r^d} \\ &\quad \times (1-t)^{d-s/2} \int_{-1}^t (1+u)^{d/2-1} (t-u)^{1-(d-s)/2-1} (1-u)^{-1} \\ &\quad \times {}_2\tilde{F}_1 \left(\begin{matrix} 1, d/2 \\ 1 - (d-s)/2 \end{matrix}; \frac{(R-1)^2}{r^2} \frac{t-u}{1-u} \right) du. \end{aligned}$$

We now apply Lemma A.1.

$$\begin{aligned} \|\epsilon_t\| &= 2^{(d-s)/2-1} \frac{\Gamma(d/2)}{\Gamma(d-s/2) \Gamma(s/2)} \frac{\omega_{d-1}}{\omega_d} \frac{(R+1)^{d-s}}{W_s(\mathbb{S}^d) r^d} (1-t)^{d/2} (1+t)^{s/2} \\ &\quad \times (1-xy)^{-d/2} \int_0^1 v^{s/2-1} (1-xv)^{d-s/2-1} \left(1 - \frac{x(1-y)}{1-xy} v \right)^{-d/2} dv, \end{aligned}$$

where $x = (1+t)/2$ and $y = (R-1)^2/r^2$. Substituting

$$1 - xy = \frac{(R+1)^2}{r^2} \frac{1-t}{2}, \quad \frac{x(1-y)}{1-xy} = \frac{4R}{(R+1)^2} \frac{1+t}{2}.$$

and (1.3) we get the Euler-type integral of an Appell function [5, Eq. 5.8(5)]

$$\begin{aligned} \|\epsilon_t\| &= \frac{2^{-s/2} \Gamma(d)}{\Gamma(d-s/2) \Gamma(s/2)} \frac{1}{W_s(\mathbb{S}^d)} (R+1)^{-s} (1+t)^{s/2} \\ &\times \int_0^1 u^{s/2-1} \left(1 - \frac{1+t}{2}u\right)^{d-s/2-1} \left(1 - \frac{4R}{(R+1)^2} \frac{1+t}{2}u\right)^{-d/2} du. \end{aligned}$$

A change of variables $1+v = (1+t)u$ yields (4.1). \square

Lemma 4.2. *Let $d-2 < s < d$. Then*

$$(4.2) \quad \|\nu_t\| = \frac{2^{1-d} \Gamma(d)}{\Gamma(d-s/2) \Gamma(s/2)} \int_{-1}^t (1+u)^{s/2-1} (1-u)^{d-s/2-1} du$$

$$(4.3) \quad = 1 - \mathbf{I}((1-t)/2; d-s/2, s/2),$$

where $\mathbf{I}(x; a, b)$ denotes the regularized incomplete Beta function (cf. (1.14)).

Proof. We proceed as in the proof of Lemma 4.1. In fact, the densities ϵ'_t and ν'_t differ by a multiplicative factor $(R+1)^{d-s}/[W_s(\mathbb{S}^d)r^d]$ and a factor $(R-1)^2/r^2$ in the argument of the hypergeometric function. From (3.2), (3.3), and Lemma A.1

$$\|\nu_t\| = \frac{\Gamma(d)}{\Gamma(d-s/2) \Gamma(s/2)} \left(\frac{1+t}{2}\right)^{s/2} \int_0^1 v^{s/2-1} \left(1 - \frac{1+t}{2}v\right)^{d-s/2-1} dv.$$

A change of variable $1+u = (1+t)v$ yields (4.2).

A manipulation of the integral (extending the integral over the complete interval $[-1, 1]$ and using the standard substitution $2v = 1-u$) yields (4.3). \square

5. THE EXTREMAL SUPPORT AND MEASURE: PROOFS OF THEOREMS 1.10, 1.11, AND 1.14.

Our first proof deals with the minimization property of S_Q .

Proof of Theorem 1.10. Let K be any compact subset of \mathbb{S}^d with positive s -capacity. For the considered range of the parameter s , we have that the potential of the extremal measure $\mu_K = \mu_{K,s}$ satisfies the following (in)equalities

$$(5.1) \quad U_s^{\mu_K}(\mathbf{x}) = W_s(K) \quad \text{q.e. on } K, \quad U_s^{\mu_K}(\mathbf{x}) \leq W_s(K) \quad \text{on } \mathbb{S}^d.$$

This follows trivially from the general theory (see [14, Chapter II]) for $d-1 \leq s < d$, with the inequality holding on the entire space \mathbb{R}^{d+1} . To derive (5.1) for the extended range, we observe that for $K = \mathbb{S}^d$ this is obvious ($\mu_K = \sigma_d$).

If $\mathbb{S}^d \setminus K$ is non-empty, there is a spherical cap Σ that contains K . The s -potential of μ_Σ equals $W_s(\Sigma)$ everywhere on Σ , so the measure $\nu := [W_s(K)/W_s(\Sigma)]\mu_K$ has a potential that equals $W_s(K)$ on Σ . Since $U_s^{\mu_K}(\mathbf{x}) \leq W_s(K)$ on $\text{supp}(\mu_K)$ (see [14, p. 136(b)]), we could derive the inequality in (5.1) by comparing the potentials of μ_K and ν and applying the restricted version of the Principle of Domination as given in [4, Lemma 5.1] (for $s = d-2$ we adapt the argument in Lemma 5.1 using [14, Theorem 1.27]). Since $U_s^{\mu_K}(\mathbf{x}) \geq W_s(K)$ q.e. on K (see [14, p. 136(a)]), we conclude the equality in (5.1) as well.

Clearly, $\mathcal{F}_s(S_Q) = F_Q$ (see (1.7) and (1.8)). We now show that for any compact set $K \subset \mathbb{S}^d$ with positive s -capacity we have $\mathcal{F}_s(K) \geq \mathcal{F}_s(S_Q)$. Indeed, let us integrate (1.7) with respect to μ_K . Since μ_K has finite energy, the inequality holds also μ_K -a.e. and we conclude that

$$\int U_s^{\mu_Q}(\mathbf{x}) \, d\mu_K(\mathbf{x}) + \int Q(\mathbf{x}) \, d\mu_K(\mathbf{x}) \geq F_Q.$$

Using the inequality in (5.1) we write

$$\int U_s^{\mu_Q}(\mathbf{x}) \, d\mu_K(\mathbf{x}) = \int U_s^{\mu_K}(\mathbf{x}) \, d\mu_Q(\mathbf{x}) \leq W_s(K),$$

which proves our claim. \square

Next, we prove sufficient conditions on Q , that guarantee that the extremal support is a spherical zone (cap).

Proof of Theorem 1.11. The convexity assumption on $f(\xi)$ implies that $Q(\mathbf{z})$ is continuous and the existence and uniqueness of the extremal measure μ_Q follows from standard potential-theoretical arguments (see [23], [24]). The rotational invariance of the external field implies that the extremal support is also rotationally invariant. Hence, there is a compact set $A \subset [-1, 1]$ and an integrable function $g : A \rightarrow \mathbb{R}^+$, such that the extremal support is given by

$$\text{supp}(\mu_Q) = \left\{ (\sqrt{1-u^2} \bar{\mathbf{x}}, u) : u \in A, \bar{\mathbf{x}} \in \mathbb{S}^{d-1} \right\},$$

and the extremal measure is

$$d\mu_Q(\mathbf{x}) = g(u) \, du \, d\sigma_{d-1}(\bar{\mathbf{x}}), \quad u \in A.$$

What we have to show is that A is connected. For this purpose we adapt the argument given in [16]. Suppose A is not connected. Then there is an interval $[\alpha, \beta] \subset (-1, 1)$, such that $[\alpha, \beta] \cap A = \{\alpha, \beta\}$. Let $A^- := A \cap [-1, \alpha]$ and $A^+ := A \cap [\beta, 1]$. Letting

$$\begin{aligned} \mathbf{x} &= (\sqrt{1-u^2} \bar{\mathbf{x}}, u), & u \in A^- \cup A^+, \bar{\mathbf{x}} \in \mathbb{S}^{d-1}, \\ \mathbf{z} &= (\sqrt{1-\xi^2} \bar{\mathbf{z}}, \xi), & \xi \in (\alpha, \beta), \bar{\mathbf{z}} \in \mathbb{S}^{d-1}, \end{aligned}$$

we represent the weighted s -potential as follows:

$$\begin{aligned} U_s^{\mu_Q}(\mathbf{z}) + Q(\mathbf{z}) &= \int_A g(u) \left(\int_{\mathbb{S}^{d-1}} \frac{d\sigma_{d-1}(\bar{\mathbf{x}})}{|\mathbf{z} - \mathbf{x}|^s} \right) du + Q(\mathbf{z}) \\ (5.2) \quad &=: \int_{A^-} g(u) \kappa(u, \xi) \, du + \int_{A^+} g(u) \kappa(u, \xi) \, du + f(\xi), \end{aligned}$$

where the kernel $\kappa(u, \xi)$ has been evaluated in [4, Section 4] for the case $\xi > u$ ($u \in A^-$) to be

$$(5.3) \quad \kappa(u, \xi) := \int_{\mathbb{S}^{d-1}} \frac{d\sigma_{d-1}(\bar{\mathbf{x}})}{|\mathbf{z} - \mathbf{x}|^s}$$

$$(5.4) \quad = (1-u)^{-s/2} (1+\xi)^{-s/2} {}_2F_1 \left(\begin{matrix} s/2, 1 - (d-s)/2 \\ d/2 \end{matrix}; \frac{1+u}{1-u} \frac{1-\xi}{1+\xi} \right)$$

$$(5.5) \quad = \sum_{k=0}^{\infty} \frac{(s/2)_k (1 - (d-s)/2)_k (1+u)^k}{(d/2)_k k! (1-u)^{k+s/2}} \frac{(1-\xi)^k}{(1+\xi)^{k+s/2}}.$$

By symmetry we derive that when $\xi < u$ ($u \in A^+$)

$$(5.6) \quad \kappa(u, \xi) = \sum_{k=0}^{\infty} \frac{(s/2)_k (1 - (d-s)/2)_k (1-u)^k}{(d/2)_k k! (1+u)^{k+s/2}} \frac{(1+\xi)^k}{(1-\xi)^{k+s/2}}.$$

It is easy to verify that the functions

$$(1-\xi)^k / (1+\xi)^{k+s/2}, \quad (1+\xi)^k / (1-\xi)^{k+s/2}, \quad k = 0, 1, 2, \dots,$$

are strictly convex for $\xi \in (-1, 1)$. Hence, from (5.5) and (5.6) we derive that the kernel $\kappa(u, \xi)$ is a convex function in ξ on (α, β) for any fixed $u \in A^- \cup A^+$. Therefore, using the convexity of $f(\xi)$ we deduce that the weighted s -potential is strictly convex on $[\alpha, \beta]$. This clearly contradicts the inequalities (1.7) and (1.8), which proves (1.26).

Now suppose that, in addition, $f(\xi)$ is also increasing. If $t_1 > -1$, for $u \in [t_1, t_2]$ and $\xi \in (-1, t_1)$, the kernel is calculated using (5.6), in which case we easily obtain that $\partial \kappa(u, \xi) / \partial \xi > 0$. This yields that the weighted s -potential is strictly increasing on $[-1, t_1]$, which contradicts (1.7) and (1.8) similarly. \square

Proof of Theorem 1.14. The external field is given by

$$Q_{\mathbf{a}, q}(\mathbf{z}) = q / |\mathbf{a} - \mathbf{z}|^s = q |R^2 - 2R\xi + 1|^{-s/2} =: f(\xi),$$

where $\mathbf{z} = (\sqrt{1 - \xi^2} \bar{\mathbf{z}}, \xi)$, $\xi \in [-1, 1]$, $\bar{\mathbf{z}} \in \mathbb{S}^{d-1}$. We easily verify that $f'(\xi) > 0$ and $f''(\xi) > 0$ for $\xi \in [-1, 1]$. According to Theorem 1.11, the extremal support associated with $Q_{\mathbf{a}, q}$ is a spherical cap. So, by Theorem 1.10 we have to minimize the \mathcal{F}_s -functional among all spherical caps centered at the South Pole.

Recall that (see (1.29) and Remark 1.13)

$$\mathcal{F}_s(\Sigma_t) = \Phi_s(t) = W_s(\mathbb{S}^d) (1 + q \|\epsilon_t\|) / \|\nu_t\|.$$

Applying the Quotient Rule and using (4.1) and (4.2) and the Fundamental Theorem of Calculus, we get (note that $\|\nu_t\| > 0$ for $t > -1$ and $\|\nu_t\|' > 0$ for $-1 < t < 1$)

$$(5.7) \quad \begin{aligned} \frac{d \Phi_s}{d t} &= \frac{q \|\epsilon_t\|' \|\nu_t\| - (1 + q \|\epsilon_t\|) \|\nu_t\|'}{\|\nu_t\|^2 / W_s(\mathbb{S}^d)} = - \frac{\|\nu_t\|'}{\|\nu_t\|} \left[\Phi_s(t) - q W_s(\mathbb{S}^d) \frac{\|\epsilon_t\|'}{\|\nu_t\|'} \right] \\ &= - \frac{\|\nu_t\|'}{\|\nu_t\|} \left[\Phi_s(t) - \frac{q (R+1)^{d-s}}{r^d} \right] =: - \frac{\|\nu_t\|'}{\|\nu_t\|} \Delta(t), \end{aligned}$$

where $r = r(t) = \sqrt{R^2 - 2Rt + 1}$. Observe, that $\Delta(t) \rightarrow \infty$ as $t \rightarrow -1$. Hence, there is a largest $t_0 \in (-1, 1]$ such that $\Delta(t) > 0$ on $(-1, t_0)$. If $t_0 = 1$, then $\Phi_s(t)$ is strictly decreasing on $(-1, 1)$ and attains its minimum at $t = 1$. We note that $\Delta(1) \geq 0$ is equivalent to the condition in Corollary 1.3. If $t_0 < 1$, then by continuity $\Delta(t_0) = 0$. Clearly, $\Phi_s'(t) < 0$ on $(-1, t_0)$ and $\Phi_s'(t_0) = 0$. Suppose, $\Phi_s'(\tau) = 0$ for some $\tau \in (-1, 1)$. Then $\Delta(\tau) = 0$. Applying the product rule we get

$$\frac{d^2 \Phi_s}{d t^2}(\tau) = - \frac{\|\nu_t\|'}{\|\nu_t\|} \left[\Phi_s'(t) - \frac{d q (R+1)^{d-s} R}{r^{d+2}} \right] \Bigg|_{t=\tau} = \frac{\|\nu_t\|'}{\|\nu_t\|} \frac{d q (R+1)^{d-s} R}{r^{d+2}} \Bigg|_{t=\tau} > 0.$$

Hence, any zero of Φ_s' is a minimum of Φ_s . Since Φ_s is twice continuously differentiable on $(-1, 1)$ (see Lemmas 4.1 and 4.2), the later observation implies that Φ_s has only one local minimum in $(-1, 1)$, namely t_0 , which has to be also a global minimum. Observe, that $\Phi_s'(t) < 0$ for $t \in (-1, t_0)$ and $\Phi_s'(t) > 0$ for $t \in (t_0, 1)$. From (5.7) we conclude that $\Delta(t) > 0$ on $(-1, t_0)$ and $\Delta(t) < 0$ on $(t_0, 1)$. This

shows that $\Phi_s(t)$ has precisely one global minimum in $(-1, 1]$, which is either the unique solution $t_0 \in (-1, 1)$ of the equation $\Delta(t) = 0$ if it exists, or $t_0 = 1$. Moreover, $\Delta(t) \geq 0$ if and only if $t \leq t_0$. By Lemma 3.3 and Remark 3.4 we have $t_0 = \max\{t : \eta_t \geq 0\}$. Clearly, $S_{Q_{a,q}} = \Sigma_{t_0}$, from the minimization property. Since the signed equilibrium for Σ_{t_0} is a positive measure, by the uniqueness of the extremal measure we derive that $\mu_{Q_{a,q}} = \eta_{t_0}$. \square

6. THE WEIGHTED s -POTENTIAL OF η_t ON $\mathbb{S}^d \setminus \Sigma_t$: ALTERNATIVE PROOF OF THEOREM 1.14

In this section we complete the proof of Theorem 1.12, namely formula (1.32) on $\mathbb{S}^d \setminus \Sigma_t$. The s -potential of η_t is given by

$$U_s^{\eta_t}(\mathbf{z}) = \int \frac{d\eta_t(\mathbf{x})}{|\mathbf{z} - \mathbf{x}|^s} = \frac{\omega_{d-1}}{\omega_d} \int_{-1}^t \kappa(u, \xi) \eta'_t(u) (1-u^2)^{d/2-1} du,$$

where $\mathbf{z} = (\sqrt{1-\xi^2} \bar{\mathbf{z}}, \xi)$, $\xi > t$, and the kernel $\kappa(u, \xi)$ is given in (5.4). The densities ϵ'_t and ν'_t of the balayage measures ϵ_t and ν_t in (1.30) have in common that they can be written as (cf. Lemmas 3.1 and 3.2)

$$(6.1) \quad \gamma'_t(u) = C \left(\frac{1-t}{1-u} \right)^{d/2} \left(\frac{t-u}{1-t} \right)^{(s-d)/2} \sum_{n=0}^{\infty} \frac{(d/2)_n}{\Gamma(n+1-(d-s)/2)} \left(c_t^2 \frac{t-u}{1-u} \right)^n$$

with appropriately chosen constants C and c_t . Hence, it is sufficient to study the s -potential of $d\gamma_t = \gamma'_t d\sigma_d|_{\Sigma_t}$.

Using the series representation (5.5) of $\kappa(u, \xi)$ and integrating term-wise we get

$$U_s^{\gamma_t}(\mathbf{z}) = \frac{C\omega_{d-1}/\omega_d}{(1+\xi)^{s/2}} \sum_{m=0}^{\infty} \sum_{n=0}^{\infty} \frac{(s/2)_m (1-(d-s)/2)_m (d/2)_n}{m! (d/2)_m \Gamma(n+1-(d-s)/2)} \left[\frac{1-\xi}{1+\xi} \right]^m c_t^{2n} \mathcal{H}_{m,n}(t; u),$$

where $\mathcal{H}_{m,n}(t; u)$ is the integral

$$\begin{aligned} \mathcal{H}_{m,n}(t; u) &= \int_{-1}^t \left(\frac{1-t}{1-u} \right)^{d/2} \left(\frac{t-u}{1-t} \right)^{(s-d)/2} \left(\frac{1+u}{1-u} \right)^m \left(\frac{t-u}{1-u} \right)^n \frac{(1-u^2)^{d/2-1}}{(1-u)^{s/2}} du \\ &= (1-t)^{d-s/2} \int_{-1}^t \frac{(t-u)^{n-(d-s)/2} (1+u)^{m+d/2-1}}{(1-u)^{m+n+1+s/2}} du. \end{aligned}$$

By [17, Eq. 2.2.6(9)]

$$\mathcal{H}_{m,n}(t; u) = \frac{\Gamma(m+d/2) \Gamma(n+1-(d-s)/2)}{\Gamma(m+n+1+s/2)} \frac{(1-t)^{d-s/2} (1+t)^{m+n+s/2}}{(1-t)^{m+d/2} (1+t)^{n+1-(d-s)/2}}.$$

Putting everything together, we arrive at

$$\begin{aligned} U_s^{\gamma_t}(\mathbf{z}) &= 2^{d-s-1} C \frac{\omega_{d-1}}{\omega_d} \frac{\Gamma(d/2)}{\Gamma(1+s/2)} \left(\frac{1-t}{2} \right)^{(d-s)/2} \left(\frac{1+t}{1+\xi} \right)^{s/2} \\ &\quad \times \sum_{m=0}^{\infty} \sum_{n=0}^{\infty} \frac{(s/2)_m (1)_n (1-(d-s)/2)_m (d/2)_n}{(1+s/2)_{m+n} m! n!} \left(\frac{1-\xi}{1+\xi} \frac{1+t}{1-t} \right)^m \left(c_t^2 \frac{1+t}{2} \right)^n. \end{aligned}$$

The double sum in the last expression is, in fact, the series expansion of the generalized F_3 -hypergeometric function (cf. [18, Eq. 7.2.4(3)])

$$F_3\left(a, a', b, b'; w, z\right) := \sum_{m=0}^{\infty} \sum_{n=0}^{\infty} \frac{(a)_m (a')_n (b)_m (b')_n}{(c)_{m+n} m! n!} w^m z^n, \quad |w|, |z| < 1.$$

Moreover, the F_3 -function in question is of the form [18, Eq. 7.2.4(76)]

$$F_3\left(a, c-a, b, c-b; w, z\right) = (1-z)^{a+b-c} {}_2F_1\left(a, b; w+z-wz\right).$$

Let r be the distance between the point charge q and any point on the boundary circle of the spherical cap Σ_t (that is $r^2 = R^2 - 2Rt + 1$) and ρ be the distance between the point charge q and \mathbf{z} on $\mathbb{S}^d \setminus \Sigma_t$ (that is $\rho^2 = |\mathbf{z} - \mathbf{a}|^2 = R^2 - 2R\xi + 1$). For $C = \Gamma(d/2)/\Gamma(d-s/2)$, $c_t = 1$ and using (1.3), we have

$$(6.2) \quad U_s^{\nu_t}(\mathbf{z}) = W_s(\mathbb{S}^d) A_{s,d} \left(\frac{1+t}{1+\xi}\right)^{s/2} {}_2F_1\left(s/2, 1 - (d-s)/2; \frac{1+t}{1+\xi}\right).$$

For $C = (1/W_s(\mathbb{S}^d)) \Gamma(d/2)/\Gamma(d-s/2)(R+1)^{d-s}/r^d$ and $c_t^2 = (R-1)^2/r^2$, we get

$$(6.3) \quad U_s^{\epsilon_t}(\mathbf{z}) = A_{s,d} \frac{1}{r^s} \left(\frac{1+t}{1+\xi}\right)^{s/2} {}_2F_1\left(s/2, 1 - (d-s)/2; \frac{\rho^2}{r^2} \frac{1+t}{1+\xi}\right).$$

The normalization constant $A_{s,d}$ is given by

$$A_{s,d} := \frac{\Gamma(d/2)}{\Gamma((d-s)/2)\Gamma(1+s/2)} = 1/{}_2F_1\left(s/2, 1 - (d-s)/2; 1\right).$$

(The above last relation holds by [1, Eq. 15.1.20].) The relations

$$\frac{1+t}{1+\xi} = \frac{(R+1)^2 - r^2}{(R+1)^2 - \rho^2}, \quad \frac{\xi-t}{1+\xi} = \frac{r^2 - \rho^2}{(R+1)^2 - \rho^2}$$

allow to express all formulas in terms of distances to the point charge q exerting the external field. Note that the hypergeometric functions above represent incomplete beta functions (see (1.13)). When using the regularized incomplete beta function $I(x; a, b)$ (see (1.14)), the s -potentials can be also written as

$$(6.4) \quad U_s^{\nu_t}(\mathbf{z}) = W_s(\mathbb{S}^d) I\left(\frac{1+t}{1+\xi}; \frac{s}{2}, \frac{d-s}{2}\right), \quad U_s^{\epsilon_t}(\mathbf{z}) = \frac{1}{\rho^s} I\left(\frac{\rho^2}{r^2} \frac{1+t}{1+\xi}; \frac{s}{2}, \frac{d-s}{2}\right),$$

which are valid for $\mathbf{z} \in \mathbb{S}^d \setminus \Sigma_t$. Hence, we obtain

$$\begin{aligned} U_s^{\eta_t}(\mathbf{z}) &= \frac{\Phi_s(t)}{W_s(\mathbb{S}^d)} U_s^{\nu_t}(\mathbf{z}) - q U_s^{\epsilon_t}(\mathbf{z}) \\ &= \Phi_s(t) I\left(\frac{1+t}{1+\xi}; \frac{s}{2}, \frac{d-s}{2}\right) - q \frac{1}{\rho^s} I\left(\frac{\rho^2}{r^2} \frac{1+t}{1+\xi}; \frac{s}{2}, \frac{d-s}{2}\right), \end{aligned}$$

By means of the functional equation $I(x; a, b) = 1 - I(1-x; b, a)$, it follows that the weighted s -potential of η_t for any $-1 < t < 1$ at \mathbf{z} in $\mathbb{S}^d \setminus \Sigma_t$ is given by

$$U_s^{\eta_t}(\mathbf{z}) + Q(\mathbf{z}) = \Phi_s(t) + \left\{ \frac{q}{\rho^s} I\left(\frac{(R+1)^2}{r^2} \frac{\xi-t}{1+\xi}; \frac{d-s}{2}, \frac{s}{2}\right) - \Phi_s(t) I\left(\frac{\xi-t}{1+\xi}; \frac{d-s}{2}, \frac{s}{2}\right) \right\},$$

which proves (1.32).

Next, we provide an alternative proof of Theorem 1.14. Using the (series) expansion

$$I(z; a, b) = \frac{\Gamma(a+b)}{\Gamma(b)} z^a (1-z)^b {}_2\tilde{F}_1\left(1, a+b; a+1; z\right),$$

we obtain for $\xi > t > -1$ the relation

$$\begin{aligned} U_s^{\eta_t}(\mathbf{z}) + Q(\mathbf{z}) &= \Phi_s(t) + \frac{\Gamma(d/2)}{\Gamma(s/2)} \left(\frac{\xi-t}{1+\xi}\right)^{(d-s)/2} \left(\frac{1+t}{1+\xi}\right)^{s/2} \\ &\times \sum_{n=0}^{\infty} \frac{(d/2)_n}{\Gamma(n+1+(d-s)/2)} \left(\frac{\xi-t}{1+\xi}\right)^n \left\{ \frac{q(R+1)^{d-s}}{r^d} \left[\frac{R^2+2R+1}{R^2-2Rt+1}\right]^n - \Phi_s(t) \right\}. \end{aligned}$$

If $q(R+1)^{d-s}/r^d \geq \Phi_s(t)$, then the above infinite series is a positive function for $1 \geq \xi > t$. An immediate consequence in such a case is the inequality

$$(6.5) \quad U_s^{\eta_t}(\mathbf{z}) + Q(\mathbf{z}) > \Phi_s(t), \quad \mathbf{z} \in \mathbb{S}^d \setminus \Sigma_t.$$

In particular, the last relation holds when $t = t_0$ is a solution of $q(R+1)^{d-s}/r^d = \Phi_s(t)$. But then from Lemma 3.3 we have that the signed equilibrium is a positive measure. Since it satisfies the Gauss variational (in)equalities (1.7) and (1.8), this is the extremal measure associated with Q . Easily, we derive that $t_0 = \max\{t : \eta_t \geq 0\}$.

Remark 6.1. An interesting observation is that for $t = t_0$ we could factor $(\xi - t)/(1 + \xi)$ (to get $[(\xi - t)/(1 + \xi)]^{1+(d-s)/2}$) and using product rule, it follows that

$$(6.6) \quad \frac{\partial}{\partial \xi} \{U_s^{\eta_t}(\mathbf{z}) + Q(\mathbf{z})\} \Big|_{\xi \rightarrow t^+} = 0.$$

It can be also shown that for $q(R+1)^{d-s}/r^d \neq \Phi_s(t)$ one has

$$\begin{aligned} \frac{\partial}{\partial \xi} \{U_s^{\eta_t}(\mathbf{z}) + Q(\mathbf{z})\} &= \frac{\Gamma(d/2)}{\Gamma((d-s)/2)\Gamma(s/2)} \left\{ \frac{q(R+1)^{d-s}}{r^d} - \Phi_s(t) \right\} \\ &\times (1+t)^{(s-d)/2} (\xi-t)^{(d-s)/2-1} + \mathcal{O}((\xi-t)^{(d-s)/2}) \quad \text{as } \xi \rightarrow t^+. \end{aligned}$$

Thus, the partial derivative with respect to ξ of the weighted s -potential of the signed equilibrium η_t is singular at the boundary of Σ_t when approaching it from the “outside” if t is not a solution of the equilibrium condition. The sign of this partial derivative is determined by the difference in curly braces, see Figure 1.

7. THE EXCEPTIONAL CASE $s = d - 2$: PROOF OF THEOREMS 1.16 AND 1.18

The proof of Theorem 1.16 will be split into several Lemmas. We first find the s -balayage of a point charge $\mathbf{y} = (\sqrt{1-v^2}\bar{\mathbf{y}}, v) \in \mathbb{S}^d \setminus \Sigma_t$ onto Σ_t . Set

$$\epsilon_{\mathbf{y}} = \epsilon_{\mathbf{y}, t, d-2} := \text{Bal}_{d-2}(\delta_{\mathbf{y}}, \Sigma_t).$$

To determine $\epsilon_{\mathbf{y}}$ we proceed as in [4, Section 3] (see also [14, Chapter IV]). We apply an inversion (stereographical projection) with center \mathbf{y} and radius $\sqrt{2}$. The image of \mathbb{S}^d is a hyperplane passing through the origin. The image of Σ_t is a hyperdisc of radius $\tau = \sqrt{1-t^2}/(v-t)$. The $(d-2)$ -extremal measure on this d -dimensional hyperdisc is the normalized (unit) uniform surface measure on its

boundary $d\lambda^*(\mathbf{x}^*) = \tau^{d-1} d\sigma_{d-1}((\mathbf{x}^* - \mathbf{b}^*)/\tau)$, where \mathbf{b}^* is the center of this hyperdisc. The potential of λ^* is found to be

$$U_{d-2}^{\lambda^*}(\mathbf{x}^*) = \tau \int_{\mathbb{S}^{d-1}} \frac{d\sigma_{d-1}((\mathbf{x}^* - \mathbf{b}^*)/\tau)}{|(\mathbf{z}^* - \mathbf{b}^*)/\tau - (\mathbf{x}^* - \mathbf{b}^*)/\tau|^{d-2}} = \tau W_{d-2}(\mathbb{S}^{d-1}) = \tau.$$

Using the Kelvin transformation of this measure as given in Section 2.1 (cf. (2.5) and (2.6) with $R^2 - 1 = 2$), we compute that

$$(7.1) \quad d\epsilon_{\mathbf{y}}(\mathbf{x}) = 2(v-t)(1-t^2)^{d/2-1} \frac{d\sigma_{d-1}(\bar{\mathbf{x}})}{|\mathbf{x} - \mathbf{y}|^d}, \quad \mathbf{x} \in \partial\Sigma_t.$$

In [4, Section 3, Eq. (3.12)] the corresponding point charge balayage was calculated for $d-2 < s < d$,

$$(7.2) \quad d\epsilon_{\mathbf{y},s}(\mathbf{x}) = \frac{2\sin(\pi(d-s)/2)}{\pi} \left(\frac{v-t}{t-u}\right)^{(d-s)/2} (1-u^2)^{d/2-1} \frac{du d\sigma_{d-1}(\bar{\mathbf{x}})}{|\mathbf{x} - \mathbf{y}|^d}, \quad \mathbf{x} \in \Sigma_t.$$

The following lemma establishes the relationship between $\epsilon_{\mathbf{y},s}$ and $\epsilon_{\mathbf{y}}$.

Lemma 7.1. *Let $d \geq 3$. Let $d\gamma_s := \frac{\sin(\pi(d-s)/2)}{\pi(t-u)^{(d-s)/2}} du$, $-1 \leq u \leq t$. Then $\|\gamma_s\| \rightarrow 1$ and $\gamma_s \xrightarrow{*} \delta_t$, as $s \rightarrow (d-2)^+$. Consequently, $\epsilon_{\mathbf{y},s} \xrightarrow{*} \epsilon_{\mathbf{y}}$, as $s \rightarrow (d-2)^+$.*

Proof. We compute

$$\|\gamma_s\| = \int_{-1}^t \frac{\sin(\pi(d-s)/2)}{\pi(t-u)^{(d-s)/2}} du = \frac{\sin(\pi(1-(d-s)/2))}{\pi(1-(d-s)/2)} (1+t)^{1-(d-s)/2}.$$

Clearly, $\|\gamma_s\| \leq 2$ and $\|\gamma_s\| \rightarrow 1$ as $s \rightarrow (d-2)^+$. Let f be a continuous function on $[-1, t]$. Then what we have to prove is that

$$\lim_{s \rightarrow (d-2)^+} \int_{-1}^t \frac{\sin(\pi(d-s)/2)}{\pi(t-u)^{(d-s)/2}} f(u) du = f(t).$$

By $\|\gamma_s\| \rightarrow 1$ as $s \rightarrow (d-2)^+$, this is equivalent to

$$(7.3) \quad \lim_{s \rightarrow (d-2)^+} \int_{-1}^t \frac{\sin(\pi(d-s)/2)}{\pi(t-u)^{(d-s)/2}} [f(u) - f(t)] du = 0.$$

Let $\epsilon > 0$. From the continuity of f it follows that there exists a $\delta > 0$ such that $|f(u) - f(t)| < \epsilon/4$ whenever $|u - t| < \delta$. For s sufficiently close to $(d-2)^+$ we estimate that

$$\left| \int_{-1}^{t-\delta} [f(u) - f(t)] \frac{\sin(\pi(d-s)/2)}{\pi(t-u)^{(d-s)/2}} du \right| \leq 2 \|f\|_{[-1,t]} \frac{\sin \pi(d-s)/2}{\pi\delta} < \epsilon/2,$$

and

$$\left| \int_{t-\delta}^t [f(u) - f(t)] \frac{\sin(\pi(d-s)/2)}{\pi(t-u)^{(d-s)/2}} du \right| < \frac{\epsilon}{4} \|\gamma_s\| \leq \epsilon/2.$$

Therefore,

$$\left| \int_{-1}^t [f(u) - f(t)] \frac{\sin(\pi(d-s)/2)}{\pi(t-u)^{(d-s)/2}} du \right| < \epsilon,$$

which proves (7.3).

Suppose now that $f(\mathbf{x})$, where $\mathbf{x} = (\sqrt{1-u^2}\bar{\mathbf{x}}, u)$, is a continuous function on \mathbb{S}^d . Then as $s \rightarrow (d-2)^+$ we have

$$\begin{aligned} \lim \int_{\Sigma_t} f \, d\epsilon_{\mathbf{y},s} &= \lim \int_{-1}^t \left(\int_{\mathbb{S}^{d-1}} f(\mathbf{x}) \frac{d\sigma_{d-1}(\bar{\mathbf{x}})}{|\mathbf{x}-\mathbf{y}|^d} \right) 2(v-t)^{(d-s)/2} (1-u^2)^{d/2-1} \, d\gamma_s(u) \\ &= 2(v-t)(1-t^2)^{d/2-1} \left(\int_{\mathbb{S}^{d-1}} f(\mathbf{x}) \frac{d\sigma_{d-1}(\bar{\mathbf{x}})}{|\mathbf{x}-\mathbf{y}|^d} \right) \Big|_{u=t} = \int_{\Sigma_t} f \, d\epsilon_{\mathbf{y}}, \end{aligned}$$

which completes the proof of the lemma. \square

Next, we determine the balayage measures in (1.34). We shall use that β_t , which is the unit charge uniformly distributed on the boundary of Σ_t , has $(d-2)$ -potential

$$(7.4) \quad U_{d-2}^{\beta_t}(\mathbf{z}) = \int_{\mathbb{S}^{d-1}, u=t} \frac{d\sigma_{d-1}(\bar{\mathbf{x}})}{|\mathbf{z}-\mathbf{x}|^{d-2}} = \begin{cases} (1-t)^{1-d/2} (1+\xi)^{1-d/2} & \text{if } \xi \geq t, \\ (1+t)^{1-d/2} (1-\xi)^{1-d/2} & \text{if } \xi < t, \end{cases}$$

where $\mathbf{z} = (\sqrt{1-\xi^2}\bar{\mathbf{z}}, u) \in \mathbb{S}^d$. This follows from (5.5) and (5.6).

Lemma 7.2. *Let $d \geq 3$. The measure $\bar{\nu}_t = \text{Bal}_{d-2}(\sigma_d, \Sigma_t)$ is given by*

$$(7.5) \quad d\bar{\nu}_t(\mathbf{x}) = d\sigma_d|_{\Sigma_t}(\mathbf{x}) + W_{d-2}(\mathbb{S}^d) \frac{1-t}{2} (1-t^2)^{d/2-1} d\delta_t(u) \, d\sigma_{d-1}(\bar{\mathbf{x}}).$$

The $(d-2)$ -potential of $\bar{\nu}_t$ is given by

$$(7.6) \quad U_{d-2}^{\bar{\nu}_t}(\mathbf{z}) = W_{d-2}(\mathbb{S}^d), \quad \mathbf{z} \in \Sigma_t,$$

$$(7.7) \quad U_{d-2}^{\bar{\nu}_t}(\mathbf{z}) = W_{d-2}(\mathbb{S}^d) (1+t)^{d/2-1} (1+\xi)^{1-d/2} < W_{d-2}(\mathbb{S}^d), \quad \mathbf{z} \in \mathbb{S}^d \setminus \Sigma_t.$$

Remark 7.3. It is interesting that the $(d-2)$ -potential of $\bar{\nu}_t$ can be expressed using the potential of β_t (cf. (7.4))

$$(7.8) \quad U_{d-2}^{\bar{\nu}_t}(\mathbf{z}) = W_{d-2}(\mathbb{S}^d) (1-t^2)^{d/2-1} U_{d-2}^{\beta_t}(\mathbf{z}), \quad \mathbf{z} \in \mathbb{S}^d \setminus \Sigma_t.$$

Remark 7.4. In the proof of Lemma 7.2 and Lemma 7.5 below we shall obtain the balayage measures constructively. Alternatively, one could get this from the potential (in)equalities (7.6), (7.7) and (7.13), (7.14).

Proof of Lemma 7.2. It is well-known that

$$(7.9) \quad \text{Bal}_{d-2}(\sigma_d, \Sigma_t) = \sigma_d|_{\Sigma_t} + \text{Bal}_{d-2}(\sigma_d|_{\mathbb{S}^d \setminus \Sigma_t}, \Sigma_t).$$

By the principle of superposition we have for $\mathbf{x} \in \partial\Sigma_t$

$$\begin{aligned} \text{Bal}_{d-2}(\sigma_d|_{\mathbb{S}^d \setminus \Sigma_t}, \Sigma_t) &= \int_{\mathbb{S}^d \setminus \Sigma_t} \epsilon_{\mathbf{y}}(\mathbf{x}) \, d\sigma_d(\mathbf{y}) \\ &= \frac{\omega_{d-1}}{\omega_d} \int_t^1 \left(\int_{\mathbb{S}^{d-1}} \epsilon_{\mathbf{y}}(\mathbf{x}) \, d\sigma_{d-1}(\bar{\mathbf{y}}) \right) (1-v^2)^{d/2-1} \, dv \\ &= 2 \frac{\omega_{d-1}}{\omega_d} (1-t^2)^{d/2-1} \left(\int_t^1 (1-v^2)^{d/2-1} (v-t) \int_{\mathbb{S}^{d-1}} \frac{d\sigma_{d-1}(\bar{\mathbf{y}})}{|\mathbf{x}-\mathbf{y}|^d} \, dv \right) \sigma_{d-1}(\bar{\mathbf{x}}). \end{aligned}$$

The inner integral can be computed using (5.6) with $s = d$

$$(7.10) \quad \int_{\mathbb{S}^{d-1}} \frac{d\sigma_{d-1}(\bar{\mathbf{y}})}{|\mathbf{x}-\mathbf{y}|^d} = \frac{1}{2(v-t)(1+v)^{d/2-1}(1-t)^{d/2-1}}.$$

Hence,

$$\begin{aligned} \text{Bal}_{d-2}(\sigma_d|_{\mathbb{S}^d \setminus \Sigma_t}, \Sigma_t) &= \frac{\omega_{d-1}}{\omega_d} (1+t)^{d/2-1} \left(\int_t^1 (1-v)^{d/2-1} dv \right) \sigma_{d-1}(\bar{\mathbf{x}}) \\ &= \frac{2}{d} \frac{\omega_{d-1}}{\omega_d} (1+t)^{d/2-1} (1-t)^{d/2} \sigma_{d-1}(\bar{\mathbf{x}}) =: q_{\bar{\mathbf{v}}_t} \sigma_{d-1}(\bar{\mathbf{x}}), \quad \mathbf{x} \in \partial \Sigma_t. \end{aligned}$$

Using $W_{d-2}(\mathbb{S}^d) = (4/d)(\omega_{d-1}/\omega_d)$ and (7.9) we derive (7.5).

Relation (7.6) holds because of the balayage properties. Using (7.4) we have

$$\begin{aligned} U_{d-2}^{\bar{\mathbf{v}}_t}(\mathbf{z}) &= \int_{\Sigma_t} \frac{d\sigma_d(\mathbf{x})}{|\mathbf{z}-\mathbf{x}|^{d-2}} + q_{\bar{\mathbf{v}}_t} U_{d-2}^{\beta_t}(\mathbf{z}) \\ &= \frac{\omega_{d-1}}{\omega_d} \int_{-1}^t (1-u^2)^{d/2-1} \int_{\mathbb{S}^{d-1}} \frac{d\sigma_{d-1}(\bar{\mathbf{x}})}{|\mathbf{z}-\mathbf{x}|^{d-2}} du + \frac{q_{\bar{\mathbf{v}}_t}}{(1-t)^{d/2-1} (1+\xi)^{d/2-1}} \\ &= \frac{\omega_{d-1}}{\omega_d} \int_{-1}^t \frac{(1-u^2)^{d/2-1}}{(1-u)^{d/2-1} (1+\xi)^{d/2-1}} du + \frac{q_{\bar{\mathbf{v}}_t}}{(1-t)^{d/2-1} (1+\xi)^{d/2-1}} \\ &= W_{d-2}(\mathbb{S}^d) \frac{1+t}{2} \frac{(1+t)^{d/2-1}}{(1+\xi)^{d/2-1}} + W_{d-2}(\mathbb{S}^d) \frac{1-t}{2} \frac{(1+t)^{d/2-1}}{(1+\xi)^{d/2-1}}, \end{aligned}$$

from which follows (7.7) □

Lemma 7.5. *Let $d \geq 3$. The measure $\bar{\epsilon}_t = \text{Bal}_{d-2}(\delta_{\mathbf{a}}, \Sigma_t)$ is given by*

$$(7.11) \quad d\bar{\epsilon}_t(\mathbf{x}) = \bar{\epsilon}'_t(u) d\sigma_d|_{\Sigma_t}(\mathbf{x}) + q_{\bar{\epsilon}_t} d\delta_t(u) d\sigma_{d-1}(\bar{\mathbf{x}}),$$

where the density $\bar{\epsilon}'_t(u)$ and the constant $q_{\bar{\epsilon}_t}$ are given by

$$(7.12) \quad \bar{\epsilon}'_t(u) := \frac{(R^2-1)^2 / W_{d-2}(\mathbb{S}^d)}{(R^2-2Ru+1)^{d/2+1}}, \quad q_{\bar{\epsilon}_t} = \frac{1-t}{2} \frac{(R+1)^2}{r^d} (1-t^2)^{d/2-1}.$$

The $(d-2)$ -potential of $\bar{\epsilon}_t$ is given by

$$(7.13) \quad U_{d-2}^{\bar{\epsilon}_t}(\mathbf{z}) = |\mathbf{z}-\mathbf{a}|^{2-d} = U_{d-2}^{\delta_{\mathbf{a}}}(\mathbf{z}), \quad \mathbf{z} \in \Sigma_t,$$

$$(7.14) \quad U_{d-2}^{\bar{\epsilon}_t}(\mathbf{z}) = r^{2-d} (1+t)^{d/2-1} (1+\xi)^{1-d/2} < U_{d-2}^{\delta_{\mathbf{a}}}(\mathbf{z}), \quad \mathbf{z} \in \mathbb{S}^d \setminus \Sigma_t.$$

Proof. As in the proof of Theorem 1.2 we evaluate

$$(7.15) \quad \bar{\epsilon}_{\mathbf{a}} := \text{Bal}_{d-2}(\delta_{\mathbf{a}}, \mathbb{S}^d), \quad d\bar{\epsilon}_{\mathbf{a}}(\mathbf{x}) = \bar{\epsilon}'_t(u) d\sigma_d(\mathbf{x}).$$

Using balayage in steps and (7.9) we get

$$(7.16) \quad \text{Bal}_{d-2}(\delta_{\mathbf{a}}, \Sigma_t) = \bar{\epsilon}_{\mathbf{a}}|_{\Sigma_t} + \text{Bal}_{d-2}(\bar{\epsilon}_{\mathbf{a}}|_{\mathbb{S}^d \setminus \Sigma_t}, \Sigma_t).$$

By the principle of superposition we have for $\mathbf{x} \in \partial \Sigma_t$

$$\begin{aligned} \text{Bal}_{d-2}(\bar{\epsilon}_{\mathbf{a}}|_{\mathbb{S}^d \setminus \Sigma_t}, \Sigma_t) &= \int_{\mathbb{S}^d \setminus \Sigma_t} \bar{\epsilon}'_t(v) \epsilon_{\mathbf{y}}(\mathbf{x}) d\sigma_d(\mathbf{y}) \\ &= \frac{\omega_{d-1}}{\omega_d} \int_t^1 \left(\int_{\mathbb{S}^{d-1}} \bar{\epsilon}'_t(v) \epsilon_{\mathbf{y}}(\mathbf{x}) d\sigma_{d-1}(\bar{\mathbf{y}}) \right) (1-v^2)^{d/2-1} dv \\ &= 2 \frac{\omega_{d-1}}{\omega_d} (1-t^2)^{d/2-1} \left(\int_t^1 (1-v^2)^{d/2-1} \bar{\epsilon}'_t(v) (v-t) \int_{\mathbb{S}^{d-1}} \frac{d\sigma_{d-1}(\bar{\mathbf{y}})}{|\mathbf{x}-\mathbf{y}|^d} dv \right) \sigma_{d-1}(\bar{\mathbf{x}}). \end{aligned}$$

Applying (7.10) yields

$$\begin{aligned} & \text{Bal}_{d-2}(\bar{\epsilon}_{\mathbf{a}}|_{\mathbb{S}^d \setminus \Sigma_t}, \Sigma_t) \\ &= \frac{\omega_{d-1} (R^2 - 1)^2}{\omega_d W_{d-2}(\mathbb{S}^d)} (1+t)^{d/2-1} \left(\int_t^1 \frac{(1-v)^{d/2-1}}{(R^2 - 2Rv + 1)^{d/2+1}} dv \right) \sigma_{d-1}(\bar{\mathbf{x}}) \\ &= \frac{2 \omega_{d-1} (R^2 - 1)^2}{d \omega_d W_{d-2}(\mathbb{S}^d)} (1+t)^{d/2-1} \frac{(1-t)^{d/2}}{(R^2 - 2Rt + 1)^{d/2}} \sigma_{d-1}(\bar{\mathbf{x}}) = q_{\bar{\epsilon}_t} \sigma_{d-1}(\bar{\mathbf{x}}), \end{aligned}$$

where we used the change of variable $w = (R-1)^2/(1-v) + 2R$ to compute the integral in the parenthesis.

Similar computations with the substitution $w = (R+1)^2/(1+u) - 2R$ (see also (7.4)) lead to (7.14). That is, for $\mathbf{z} \in \mathbb{S}^d \setminus \Sigma_t$ one has

$$\begin{aligned} U_{d-2}^{\bar{\epsilon}_t}(\mathbf{z}) &= \int_{\Sigma_t} \frac{\bar{\epsilon}'_t(u) d\sigma_d(\mathbf{x})}{|\mathbf{z} - \mathbf{x}|^{d-2}} + q_{\bar{\epsilon}_t} U_{d-2}^{\beta_t}(\mathbf{z}) \\ &= \frac{\omega_{d-1} (R^2 - 1)^2}{\omega_d W_{d-2}(\mathbb{S}^d)} \int_{-1}^t \frac{(1-u^2)^{d/2-1}}{(R^2 - 2Ru + 1)^{d/2+1}} \int_{\mathbb{S}^{d-1}} \frac{d\sigma_{d-1}(\bar{\mathbf{x}})}{|\mathbf{z} - \mathbf{x}|^{d-2}} du + q_{\bar{\epsilon}_t} U_{d-2}^{\beta_t}(\mathbf{z}) \\ &= \frac{\omega_{d-1} (R^2 - 1)^2}{\omega_d W_{d-2}(\mathbb{S}^d)} \int_{-1}^t \frac{(1+u)^{d/2-1}}{(R^2 - 2Ru + 1)^{d/2+1} (1+\xi)^{d/2-1}} du + q_{\bar{\epsilon}_t} U_{d-2}^{\beta_t}(\mathbf{z}) \\ &= \frac{2 \omega_{d-1} (R-1)^2}{d \omega_d W_{d-2}(\mathbb{S}^d) r^d} \frac{(1+t)^{d/2}}{(1+\xi)^{d/2-1}} + \frac{1-t}{2} \frac{(R+1)^2}{r^d} \frac{(1+t)^{d/2-1}}{(1+\xi)^{d/2-1}} \\ &= \frac{1}{r^{d-2}} \frac{(1+t)^{d/2-1}}{(1+\xi)^{d/2-1}} \left[\frac{(R-1)^2}{R^2 - 2Rt + 1} \frac{1+t}{2} + \frac{1-t}{2} \frac{(R+1)^2}{R^2 - 2Rt + 1} \right] = \frac{1}{r^{d-2}} \frac{(1+t)^{d/2-1}}{(1+\xi)^{d/2-1}}. \end{aligned}$$

As in the proof of Lemma 7.2 the balayage properties imply Equation (7.13). \square

The weak* convergence in (1.36) is shown next.

Lemma 7.6. *Let $t \in (-1, 1)$ be fixed. Then*

$$(7.17) \quad \nu_{t,s} \xrightarrow{*} \bar{\nu}_t, \quad \epsilon_{t,s} \xrightarrow{*} \bar{\epsilon}_t, \quad \text{as } s \rightarrow (d-2)^+.$$

Proof. The result follows easily from the weak* convergence $\epsilon_{\mathbf{y},s} \xrightarrow{*} \epsilon_{\mathbf{y}}$ as $s \rightarrow (d-2)^+$ and the following representation valid for any measure μ on \mathbb{S}^d :

$$(7.18) \quad \text{Bal}_s(\mu, \Sigma_t)(\mathbf{x}) = \mu|_{\Sigma_t}(\mathbf{x}) + \int_{\mathbb{S}^d \setminus \Sigma_t} \epsilon_{\mathbf{y},s}(\mathbf{x}) d\mu(\mathbf{y}).$$

\square

The norms $\|\bar{\nu}_t\|$ and $\|\bar{\epsilon}_t\|$ can be obtained from Lemmas 4.1 and 4.2 by taking the limit $s \rightarrow (d-2)^+$ (which is justified by the weak* convergence shown in Lemma 7.1).

Lemma 7.7. *Let $d \geq 3$. Then*

$$(7.19) \quad \|\bar{\epsilon}_t\| = \frac{d-2}{4} (R+1)^2 \int_{-1}^t \frac{(1+u)^{d/2-2} (1-u)^{d/2}}{(R^2 - 2Ru + 1)^{d/2}} du,$$

$$(7.20) \quad \|\bar{\nu}_t\| = \frac{d-2}{4} W_{d-2}(\mathbb{S}^d) \int_{-1}^t (1+u)^{d/2-2} (1-u)^{d/2} du.$$

Completion of the proof of Theorem 1.16. Proceeding as in the proof of Theorem 1.14, but using now $(r = r(t) = \sqrt{R^2 - 2Rt + 1})$

$$\overline{\Phi}'_{d-2}(t) = -\|\overline{v}_t\|' / \|\overline{v}_t\| \left[\overline{\Phi}_{d-2}(t) - q(R+1)^2 / r^d \right] =: -\|\overline{v}_t\|' / \|\overline{v}_t\| \Delta(t),$$

it follows that the global minimum of $\overline{\Phi}_{d-2}$ is either the unique solution $t_0 \in (-1, 1)$ of the equation $\Delta(t) = 0$, or $t_0 = 1$. In particular, $\Delta(t) \geq 0$ if and only if $t \leq t_0$.

The explicite form (1.35) follows from Lemmas 7.2 and 7.5. If $\overline{\eta}_t \geq 0$ then $\Delta(t) \geq 0$, so $t \leq t_0$. On the other hand, it is easy to see that if $t = t_0$, then $\overline{\eta}_{t_0}$ given in (1.38) is ≥ 0 because of $(R-1)^2 < R^2 - 2Rt_0 + 1 < R^2 - 2Ru + 1$. Therefore, we have that $t_0 = \max\{t : \overline{\eta}_t \geq 0\}$, $\mu_{\overline{Q}_{\mathbf{a},q}} = \overline{\eta}_{t_0}$, and $\text{supp}(\mu_{\overline{Q}_{\mathbf{a},q}}) = \Sigma_{t_0}$. \square

The proof of Theorem 1.18 is also split into several lemmas.

We must check that Theorem 1.11 also holds in the case $d = 2$ and $s = 0$. Then we can make use of the fact that the support $S_{\overline{Q}_{\mathbf{a},q}}$ of the extremal measure on \mathbb{S}^2 associated with the external logarithmic field $\overline{Q}_{\mathbf{a},q}$ is a spherical cap.

Adaptation of the proof of Theorem 1.11 for $d = 2$ and $s = 0$. Theorem 1.11 can be extended to hold for $d = 2$ and $s = 0$. Instead of the kernel $\kappa(u, \xi)$ given in (5.3) one has to consider

$$\kappa_0(u, \xi) := \int_{\mathbb{S}^1} \log \frac{1}{|\mathbf{z} - \mathbf{x}|} d\sigma_1(\overline{\mathbf{x}}) = -\frac{1}{2} \frac{1}{\pi} \int_{-1}^1 \frac{\log \left(2 - 2u\xi - 2\sqrt{1-u^2}\sqrt{1-\xi^2}\tau \right)}{\sqrt{1-\tau^2}} d\tau \quad (7.21)$$

$$= -\frac{1}{2} \log(1 - u\xi + |\xi - u|) = \begin{cases} -\frac{1}{2} \log(1 + \xi) - \frac{1}{2} \log(1 - u) & \xi \geq u, \\ -\frac{1}{2} \log(1 - \xi) - \frac{1}{2} \log(1 + u) & \xi \leq u. \end{cases}$$

This follows from the Funk-Hecke formula and [21, Lemma 1.15]. It is easy to verify that the kernel $\kappa_0(u, \xi)$ is strictly convex for $\xi \in (-1, 1)$ for any fixed $u \in (-1, 1)$. Hence, we may use the arguments of the proof of Theorem 1.11 appropriately adapted for $d = 2$ and $s = 0$. \square

It should be emphasized that in the logarithmic case balayage preserves mass. Thus, the logarithmic potentials of a measure and its logarithmic balayage onto a compact set K differ by a constant on K .

Lemma 7.8. *Let $d = 2$ and $s = 0$. The measure $\overline{v}_{t,0} = \text{Bal}_0(\sigma_2, \Sigma_t)$ is given by*

$$(7.22) \quad d\overline{v}_{t,0}(\mathbf{x}) = d\sigma_2|_{\Sigma_t}(\mathbf{x}) + \frac{1-t}{2} d\delta_t(u) d\sigma_1(\overline{\mathbf{x}})$$

and $\|\overline{v}_{t,0}(\mathbf{x})\| = 1$. The logarithmic potential of $\overline{v}_{t,0}$ is given by

$$U_0^{\overline{v}_{t,0}}(\mathbf{z}) = \frac{1+t}{4} - \frac{\log 2}{2} - \frac{1}{2} \log(1+t), \quad \mathbf{z} \in \Sigma_t,$$

$$U_0^{\overline{v}_{t,0}}(\mathbf{z}) = \frac{1+t}{4} - \frac{\log 2}{2} - \frac{1}{2} \log(1+\xi), \quad \mathbf{z} \in \mathbb{S}^d \setminus \Sigma_t.$$

The measure $\overline{v}_{t,0}$ is the logarithmic extremal measure on Σ_t and

$$(7.23) \quad W_0(\Sigma_t) = \frac{1+t}{4} - \frac{\log 2}{2} - \frac{1}{2} \log(1+t).$$

Proof. Using relation (7.21) we show that the measure in (7.22) satisfies the balayage properties. Let $\mathbf{z} \in \Sigma_t$, that is $\xi \leq t$. Then

$$\begin{aligned} U_0^{\bar{\nu}_{t,0}}(\mathbf{z}) &= U_0^{\sigma_2}(\mathbf{z}) - U_0^{\sigma_2|_{\mathbb{S}^2 \setminus \Sigma_t}}(\mathbf{z}) + \frac{1-t}{2} U_0^{\sigma_1|_{u=t}}(\mathbf{z}) \\ &= W_0(\mathbb{S}^2) - \frac{\omega_1}{\omega_2} \int_t^1 \left(\int_{\mathbb{S}^1} \log \frac{1}{|\mathbf{z} - \mathbf{x}|} d\sigma_1(\bar{\mathbf{x}}) \right) du + \frac{1-t}{2} \int_{\mathbb{S}^1} \log \frac{1}{|\mathbf{z} - \mathbf{x}|} \Big|_{u=t} d\sigma_1(\bar{\mathbf{x}}) \\ &= \frac{1}{2} - \log 2 + \frac{1}{4} \int_{-1}^t [\log(1-\xi) + \log(1+u)] du - \frac{1-t}{4} [\log(1-\xi) + \log(1+t)] \\ &= \frac{1+t}{4} - \frac{\log 2}{2} - \frac{1}{2} \log(1+t) = W_0(\Sigma_t). \end{aligned}$$

For $\mathbf{z} \in \mathbb{S}^2 \setminus \Sigma_t$, that is $\xi \geq t \geq u$, we have after a similar computation

$$\begin{aligned} U_0^{\bar{\nu}_{t,0}}(\mathbf{z}) &= \int_{\Sigma_t} \log \frac{1}{|\mathbf{z} - \mathbf{x}|} d\sigma_2(\mathbf{x}) + \frac{1-t}{2} \int_{\mathbb{S}^1} \log \frac{1}{|\mathbf{z} - \mathbf{x}|} \Big|_{u=t} d\sigma_1(\bar{\mathbf{x}}) \\ &= \frac{\omega_1}{\omega_2} \int_{-1}^t \left(\int_{\mathbb{S}^1} \log \frac{1}{|\mathbf{z} - \mathbf{x}|} d\sigma_1(\bar{\mathbf{x}}) \right) du + \frac{1-t}{2} \int_{\mathbb{S}^1} \log \frac{1}{|\mathbf{z} - \mathbf{x}|} \Big|_{u=t} d\sigma_1(\bar{\mathbf{x}}) \\ &= \frac{1+t}{4} - \frac{\log 2}{2} - \frac{1}{2} \log(1+\xi) = W_0(\Sigma_t) + \frac{1}{2} \log \frac{1+t}{1+\xi} < W_0(\Sigma_t). \end{aligned}$$

Since

$$\|\nu_{t,0}\| = \int_{\Sigma_t} d\sigma_2 + \frac{1-t}{2} \int_{\mathbb{S}^1} d\sigma_1 = \frac{\omega_1}{\omega_2} \int_{-1}^t du + \frac{1-t}{2} \int_{\mathbb{S}^1} d\sigma_1 = \frac{1+t}{2} + \frac{1-t}{2} = 1,$$

$\nu_{t,0}$ is a probability measure on Σ_t which is constant there. By uniqueness of the logarithmic extremal measure μ_{Σ_t} on Σ_t one has $\mu_{\Sigma_t} = \nu_{t,0}$. \square

Lemma 7.9. *Let $d = 2$ and $s = 0$. Then the Mhaskar-Saff functional \mathcal{F}_0 for spherical caps Σ_t is given by*

$$(7.24) \quad \begin{aligned} \mathcal{F}_0(\Sigma_t) &= (1+q) \frac{1+t}{4} + q \frac{(R-1)^2 \log(R^2 - 2Rt + 1)}{8R} - \frac{1}{2} \log(1+t) \\ &\quad - \frac{\log 2}{2} - q \frac{(R+1)^2 \log(R+1)^2}{8R}. \end{aligned}$$

It has precisely one global minimum $t_0 \in (-1, 1]$. This minimum is given by

$$(7.25) \quad t_0 = \min \left\{ 1, \frac{R^2 - 2Rq + 1}{[2R(1+q)]} \right\}.$$

Proof. By Lemma 7.8 and $|\mathbf{x} - \mathbf{a}|^2 = R^2 - 2Ru + 1$ we obtain (with $\mu_{\Sigma_t,0} = \bar{\nu}_{t,0}$)

$$\begin{aligned} \int \bar{Q}_{\mathbf{a},q} d\mu_{\Sigma_t,0} &= q \int_{\Sigma_t} \log \frac{1}{|\mathbf{x} - \mathbf{a}|} d\sigma_2(\mathbf{x}) + q \frac{1-t}{2} \int_{\mathbb{S}^1} \log \frac{1}{|\mathbf{x} - \mathbf{a}|} \Big|_{u=t} d\sigma_1(\bar{\mathbf{x}}) \\ &= -\frac{q\omega_1}{2\omega_2} \int_{-1}^t \log(R^2 - 2Ru + 1) du - \frac{q}{2} \frac{1-t}{2} \log(R^2 - 2Rt + 1) \\ &= q \frac{1+t}{4} - q \frac{(R+1)^2 \log(R+1)^2}{8R} + q \frac{(R-1)^2 \log(R^2 - 2Rt + 1)}{8R}. \end{aligned}$$

Substitution of the last expression and $W_0(\Sigma_t)$ from (7.23) into

$$\mathcal{F}_0(t) := \mathcal{F}_0(\Sigma_t) = W_0(\Sigma_t) + \int \bar{Q}_{\mathbf{a},q} d\mu_{\Sigma_t,0},$$

yields (7.24). Observe, that $\mathcal{F}_0(t) \rightarrow \infty$ as $t \rightarrow -1$. Furthermore,

$$\begin{aligned} \mathcal{F}'_0(t) &= \frac{1+q}{4} - q \frac{(R-1)^2}{4(R^2-2Rt+1)} - \frac{1}{2(1+t)} = \frac{1-t}{4} \left(\frac{2qR}{R^2-2Rt+1} - \frac{1}{1+t} \right) \\ &= \frac{R(1+q)(1-t)}{2(1+t)(R^2-2Rt+1)} \left[1+t - \frac{(R+1)^2}{2R(1+q)} \right]. \end{aligned}$$

If $-1 < t < 1$, then the sign of $\mathcal{F}'_0(t)$ is given by the sign of the linear function in the brackets, which is negative at $t = -1$. If $(R+1)^2 \geq 4R(1+q)$, then $\mathcal{F}'_0(t) < 0$ everywhere on $(-1, 1)$, and $\mathcal{F}_0(\Sigma_t)$ is strictly monotonically decreasing on $(-1, 1)$ and has a global minimum at $t = 1$. Otherwise, if $(R+1)^2 < 4R(1+q)$, then $\mathcal{F}'_0(t)$ has exactly one zero $t_0 := (R^2 - 2Rq + 1) / [2R(1+q)]$ on $(-1, 1)$, and is negative on $(-1, t_0)$ and positive on $(t_0, 1)$. Clearly, $\mathcal{F}_0(t)$ achieves global minimum on $(-1, 1]$ at t_0 , with value

$$\mathcal{F}_0(\Sigma_{t_0}) = \frac{(R+1)^2}{8R} + q \frac{(R-1)^2}{8R} \log \frac{q}{1+q} - \frac{1}{2} \log \frac{(R+1)^2}{R(1+q)} - q \log(R+1).$$

This completes the proof. \square

Lemma 7.10. *Let $d = 2$ and $s = 0$. The measure $\bar{\epsilon}_{t,0} = \text{Bal}_0(\delta_{\mathbf{a}}, \Sigma_t)$ is given by*

$$(7.26) \quad d\bar{\epsilon}_{t,0}(\mathbf{x}) = \frac{(R^2-1)^2}{(R^2-2Ru+1)^2} d\sigma_2|_{\Sigma_t}(\mathbf{x}) + \frac{1-t}{2} \frac{(R+1)^2}{R^2-2Rt+1} d\delta_t(u) d\sigma_1(\bar{\mathbf{x}})$$

and $\|\bar{\epsilon}_{t,0}\| = 1$. The logarithmic potential of $\bar{\epsilon}_{t,0}$ is given by

$$\begin{aligned} U_0^{\bar{\epsilon}_{t,0}}(\mathbf{z}) &= U_0^{\delta_{\mathbf{a}}}(\mathbf{z}) + \frac{1}{2} \log \frac{R^2-2Rt+1}{2(1+t)} + \frac{(R+1)^2}{8R} \log \frac{(R+1)^2}{R^2-2Rt+1}, \quad \mathbf{z} \in \Sigma_t, \\ U_{d-2}^{\bar{\epsilon}_{t,0}}(\mathbf{z}) &= U_0^{\delta_{\mathbf{a}}}(\mathbf{z}) + \frac{1}{2} \log \frac{R^2-2R\xi+1}{2(1+\xi)} + \frac{(R+1)^2}{8R} \log \frac{(R+1)^2}{R^2-2Rt+1}, \quad \mathbf{z} \in \mathbb{S}^d \setminus \Sigma_t. \end{aligned}$$

Proof. Let $\mathbf{z} \in \Sigma_t$. We write

$$\begin{aligned} U_0^{\bar{\epsilon}_{t,0}}(\mathbf{z}) &= \frac{\omega_1}{\omega_2} \left(\int_{-1}^{\xi} + \int_{\xi}^t \right) \frac{(R^2-1)^2}{(R^2-2Ru+1)^2} \left(\int_{\mathbb{S}^1} \log \frac{1}{|\mathbf{z}-\bar{\mathbf{x}}|} d\sigma_1(\bar{\mathbf{x}}) \right) du \\ &\quad + \frac{1-t}{2} \frac{(R+1)^2}{R^2-2Rt+1} \int_{\mathbb{S}^1} \log \frac{1}{|\mathbf{z}-\bar{\mathbf{x}}|} \Big|_{u=t} d\sigma_1(\bar{\mathbf{x}}). \end{aligned}$$

Using relation (7.21) and Mathematica we arrive at

$$(7.27) \quad U_0^{\bar{\epsilon}_{t,0}}(\mathbf{z}) = -\frac{1}{2} \log(R^2-2R\xi+1) + C(R; t),$$

where

$$C(R, t) := \frac{1}{2} \log \frac{R^2-2Rt+1}{2(1+t)} + \frac{(R+1)^2}{8R} \log \frac{(R+1)^2}{R^2-2Rt+1}.$$

Let $\mathbf{z} \in \mathbb{S}^2 \setminus \Sigma_t$. Then

$$\begin{aligned} U_0^{\bar{\epsilon}_{t,0}}(\mathbf{z}) &= \frac{\omega_1}{\omega_2} \int_{-1}^t \frac{(R^2-1)^2}{(R^2-2Ru+1)^2} \left(\int_{\mathbb{S}^1} \log \frac{1}{|\mathbf{z}-\bar{\mathbf{x}}|} d\sigma_1(\bar{\mathbf{x}}) \right) du \\ &\quad + \frac{1-t}{2} \frac{(R+1)^2}{R^2-2Rt+1} \int_{\mathbb{S}^1} \log \frac{1}{|\mathbf{z}-\bar{\mathbf{x}}|} \Big|_{u=t} d\sigma_1(\bar{\mathbf{x}}) \end{aligned}$$

Using relation (7.21) and evaluating the integral one gets after some simplifications

$$(7.28) \quad U_0^{\bar{\epsilon}_{t,0}}(\mathbf{z}) = -\frac{1}{2} \log [2(1+\xi)] + \frac{(R+1)^2}{8R} \log \frac{(R+1)^2}{R^2 - 2Rt + 1},$$

which yields the representation outside of Σ_t . Since

$$\frac{R^2 - 2R\xi + 1}{R^2 - 2Rt + 1} \frac{1+t}{1+\xi} < 1 \quad \text{for } \xi > t,$$

it follows for $\mathbf{z} \in \mathbb{S}^2 \setminus \Sigma_t$ that

$$U_0^{\bar{\epsilon}_{t,0}}(\mathbf{z}) = U_0^{\delta_{\mathbf{a}}}(\mathbf{z}) + C(R, t) + \frac{1}{2} \log \left[\frac{R^2 - 2R\xi + 1}{R^2 - 2Rt + 1} \frac{1+t}{1+\xi} \right] < U_0^{\delta_{\mathbf{a}}}(\mathbf{z}) + C(R, t).$$

Hence, $\bar{\epsilon}_{t,0}$ has the properties of a logarithmic balayage measure. Finally,

$$\begin{aligned} \|\bar{\epsilon}_{t,0}\| &= \int_{\Sigma_t} \frac{(R^2 - 1)^2}{(R^2 - 2Ru + 1)^2} d\sigma_2(\mathbf{x}) + \frac{1-t}{2} \frac{(R+1)^2}{R^2 - 2Rt + 1} \int_{\mathbb{S}^1} d\sigma_1(\bar{\mathbf{x}}) \\ &= \frac{\omega_1}{\omega_2} \int_{-1}^t \frac{(R^2 - 1)^2}{(R^2 - 2Ru + 1)^2} du + \frac{1-t}{2} \frac{(R+1)^2}{R^2 - 2Rt + 1} \\ &= \frac{1+t}{2} \frac{(R-1)^2}{R^2 - 2Rt + 1} + \frac{1-t}{2} \frac{(R+1)^2}{R^2 - 2Rt + 1} = 1. \end{aligned}$$

This completes the proof. \square

Proof of Theorem 1.18. Lemmas 7.8 and 7.10 imply that $\bar{\eta}_{t,0} = (1+q)\bar{\nu}_{t,0} - q\bar{\epsilon}_{t,0}$ is, indeed, the logarithmic signed equilibrium on Σ_t associated with $\bar{Q}_{\mathbf{a},q}$ as can be seen from its weighted logarithmic potential given in the Theorem. Using $r = \sqrt{R^2 - 2Rt + 1}$ and $\rho = \sqrt{R^2 - 2Ru + 1}$, we can write

$$d\bar{\eta}_{t,0}(\mathbf{x}) = \left[1 + q - \frac{q(R^2 - 1)^2}{\rho^4} \right] d\sigma_2|_{\Sigma_t}(\mathbf{x}) + \frac{1-t}{2} \left[1 + q - \frac{q(R+1)^2}{r^2} \right] d\beta_t(\mathbf{x}),$$

where $\mathbf{x} \in \Sigma_t$. If $\bar{\eta}_{t,0} \geq 0$, then $1+q - q(R+1)^2/(R^2 - 2Rt + 1) \geq 0$, so $t \leq t_0$. On the other hand, it is easy to see that if $t = t_0$, then $\bar{\eta}_{t_0,0}$ given in (1.41) is ≥ 0 because $\rho \leq \rho$ and $(R-1)^2 < R^2 - 2Ru + 1$. Therefore, we have that $t_0 = \max\{t : \bar{\eta}_{t,0} \geq 0\}$, $\mu_{\bar{Q}_{\mathbf{a},q}} = \bar{\eta}_{t_0,0}$, and $\text{supp}(\mu_{\bar{Q}_{\mathbf{a},q}}) = \Sigma_{t_0}$. \square

8. AXIS-SUPPORTED RIESZ EXTERNAL FIELDS

In this section we shall prove Theorems 1.23, 1.24, 1.25, 1.26, and 1.27.

Proof of Theorem 1.23. Direct calculation shows that

$$\begin{aligned} U_s^{\bar{\eta}^\lambda}(\mathbf{z}) &= \frac{\mathcal{F}_s(\mathbb{S}^d)}{W_s(\mathbb{S}^d)} U_s^{\sigma_d}(\mathbf{z}) - \int \left(\int_{\mathbb{S}^d} \frac{(R^2 - 1)^{d-s}}{|\mathbf{z} - \mathbf{x}|^s |\mathbf{x} - \mathbf{a}|^{2d-s}} d\sigma_d(\mathbf{x}) \right) d\lambda(R) \\ &= \frac{\mathcal{F}_s(\mathbb{S}^d)}{W_s(\mathbb{S}^d)} W_s(\mathbb{S}^d) - \int \frac{d\lambda(R)}{|\mathbf{z} - R\mathbf{p}|^s} = \mathcal{F}_s(\mathbb{S}^d) - Q(\mathbf{z}), \end{aligned}$$

where we used the Kelvin transformation for points (cf. proof of Theorem 1.2).

The second part follows from the uniqueness of the s -extremal measure on \mathbb{S}^d associated with Q and the fact that the density is minimal at the North Pole. \square

Proof of Theorem 1.24. The logarithmic potential of $\tilde{\eta}_{\lambda,0}$ is given by

$$U_0^{\tilde{\eta}_{\lambda,0}}(\mathbf{z}) = (1 + \|\lambda\|) W_0(\mathbb{S}^d) - \int \left(\int_{\mathbb{S}^d} \frac{(R^2 - 1)^d}{|\mathbf{x} - R\mathbf{p}|^{2d}} \log \frac{1}{|\mathbf{z} - \mathbf{x}|} d\sigma(\mathbf{x}) \right) d\lambda(R).$$

A Kelvin transformation with center $\mathbf{a} = R\mathbf{p}$ and radius $\sqrt{R^2 - 1}$ (cf. Section 2.1) yields

$$\begin{aligned} \int_{\mathbb{S}^d} \frac{(R^2 - 1)^d}{|\mathbf{x} - R\mathbf{p}|^{2d}} \log \frac{1}{|\mathbf{z} - \mathbf{x}|} d\sigma(\mathbf{x}) &= \int_{\mathbb{S}^d} \frac{(R^2 - 1)^d}{|\mathbf{x} - \mathbf{a}|^d} \log \frac{1}{|\mathbf{z} - \mathbf{x}|} \frac{d\sigma(\mathbf{x})}{|\mathbf{x} - \mathbf{a}|^d} \\ &= \int_{(\mathbb{S}^d)^*} |\mathbf{x}^* - \mathbf{a}|^d \log \frac{|\mathbf{z}^* - \mathbf{a}|^d |\mathbf{x}^* - \mathbf{a}|^d}{(R^2 - 1) |\mathbf{z}^* - \mathbf{x}^*|^d |\mathbf{x}^* - \mathbf{a}|^d} d\sigma(\mathbf{x}^*) \\ &= \log \frac{|\mathbf{z}^* - \mathbf{a}|}{R^2 - 1} - \int_{(\mathbb{S}^d)^*} \log \frac{1}{|\mathbf{x}^* - \mathbf{a}|} d\sigma(\mathbf{x}^*) + \int_{(\mathbb{S}^d)^*} \log \frac{1}{|\mathbf{z}^* - \mathbf{x}|} d\sigma(\mathbf{x}^*) \\ &= \log \frac{1}{|\mathbf{z} - \mathbf{a}|} - \int_{\mathbb{S}^d} \log \frac{1}{|\mathbf{y} - \mathbf{a}|} d\sigma(\mathbf{y}) + U_0^\sigma(\mathbf{z}). \end{aligned}$$

Hence

$$\begin{aligned} U_0^{\tilde{\eta}_{\lambda,0}}(\mathbf{z}) &= (1 + \|\lambda\|) W_0(\mathbb{S}^d) - \int \log \frac{1}{|\mathbf{z} - \mathbf{a}|} d\lambda(R) \\ &\quad + \int_{\mathbb{S}^d} \left(\int \log \frac{1}{|\mathbf{y} - \mathbf{a}|} d\lambda(R) \right) d\sigma(\mathbf{y}) - \|\lambda\| W_0(\mathbb{S}^d), \end{aligned}$$

from which follows the first part of the theorem.

The second part follows from the uniqueness of the logarithmic extremal measure on \mathbb{S}^d associated with Q (Lemma 2.1 and in particular [19]) and the fact that the density is minimal at the North Pole. \square

Proof of Theorem 1.25. By construction $\tilde{\eta}_t$ is of total charge one. It is easy to verify that the signed measure (1.54) has a constant weighted s -potential on Σ_t . Indeed,

$$U_s^{\tilde{\epsilon}_t}(\mathbf{x}) = \int U_s^{\text{Bal}_s(\delta_{R\mathbf{p}}, \Sigma_t)}(\mathbf{x}) d\lambda(R) = \int \frac{d\lambda(R)}{|\mathbf{x} - R\mathbf{p}|^s} = Q(\mathbf{x}), \quad \mathbf{x} \in \Sigma_t.$$

Together with $U_s^{\nu_t}(\mathbf{z}) = W_s(\mathbb{S}^d)$ on Σ_t we have $U_s^{\tilde{\eta}_t}(\mathbf{z}) = \tilde{\Phi}_s(t)$ on Σ_t . Moreover, by Remark 1.8, we also have that $\mathcal{F}_s(\Sigma_t) = \tilde{\Phi}_s(t)$.

By definition of ν_t , $\tilde{\epsilon}_t$, and $\text{Bal}_s(\delta_{R\mathbf{p}}, \Sigma_t) = \epsilon_{t,R}$ (with additional indication of the dependence on the parameter R) we can write

$$\tilde{\eta}_t = \frac{\tilde{\Phi}_s(t)}{W_s(\mathbb{S}^d) \|\lambda\|} \int \nu_t d\lambda(R) - \int \epsilon_{t,R} d\lambda(R) = \int \left[\frac{\tilde{\Phi}_s(t)}{W_s(\mathbb{S}^d) \|\lambda\|} \nu_t - \epsilon_{t,R} \right] d\lambda(R).$$

Thus, the signed equilibrium is

$$d\tilde{\eta}_t(\mathbf{x}) = \left[\int \tilde{\eta}_t''(u, R) d\lambda(R) \right] \frac{\omega_{d-1}}{\omega_d} (1 - u^2)^{d/2-1} du d\sigma_{d-1}(\bar{\mathbf{x}}), \quad \mathbf{x} \in \Sigma_t,$$

where, when using Lemmas 3.1 and 3.2, we have

$$(8.1) \quad \begin{aligned} \tilde{\eta}_t''(u, R) = & \frac{1}{W_s(\mathbb{S}^d)} \frac{1}{\|\lambda\|} \frac{\Gamma(d/2)}{\Gamma(d-s/2)} \left(\frac{1-t}{1-u} \right)^{d/2} \left(\frac{t-u}{1-t} \right)^{(s-d)/2} \\ & \times \left\{ \tilde{\Phi}_s(t) {}_2\tilde{F}_1 \left(\begin{matrix} 1, d/2 \\ 1 - (d-s)/2 \end{matrix}; \frac{t-u}{1-u} \right) \right. \\ & \left. - \frac{\|\lambda\| (R+1)^{d-s}}{r^d} {}_2\tilde{F}_1 \left(\begin{matrix} 1, d/2 \\ 1 - (d-s)/2 \end{matrix}; \frac{(R-1)^2}{r^2} \frac{t-u}{1-u} \right) \right\}. \end{aligned}$$

We claim that the density (the integral in square brackets) is either positive for all $u \in [-1, t]$, or is positive on some interval $[-1, t_c)$ and negative on $(t_c, t]$. It suffices to consider the function $h(u)$ obtained by integrating the expression in braces in (8.1) against $d\lambda(R)$. Using the series expansion of the hypergeometric functions we get

$$(8.2) \quad \begin{aligned} h(u) = & \sum_{k=0}^{\infty} \frac{(d/2)_k}{\Gamma(k+1 - (d-s)/2)} \left(\frac{t-u}{1-u} \right)^k \\ & \times \left\{ \int \left[\tilde{\Phi}_s(t) - \frac{\|\lambda\| (R+1)^{d-s}}{r^d} \left(\frac{R-1}{r} \right)^{2k} \right] d\lambda(R) \right\}. \end{aligned}$$

The coefficients in braces form an increasing sequence with positive limit as $k \rightarrow \infty$. Hence, either all coefficients are positive, or the first n are negative and then all others are positive. So, with a substitution $x = (t-u)/(1-u)$ ($x \in A_t := [0, (1+t)/2]$) we obtain

$$g(x) = \sum_{k=0}^{\infty} \frac{a_k}{k!} x^k, \quad a_k < 0 \text{ for } k < n \text{ and } a_k \geq 0 \text{ for } k \geq n.$$

We have that $g^{(n)}(x) > 0$ on A_t , so $g^{(n-1)}(x)$ is strictly increasing on A_t . Since $g^{(n-1)}(0) = a_{n-1} < 0$, there is a γ_{n-1} in A_t such that $g^{(n-1)}(x)$ is negative on $[0, \gamma_{n-1})$ and positive on $(\gamma_{n-1}, (1+t)/2]$. Indeed, if such a γ_{n-1} does not exist, we get a contradiction, because $g^{(n-1)}(x)$ will be negative on A_t , which would imply that $g^{(n-2)}(x)$ is decreasing and negative on A_t , and so on. This argument yields $g(x) < 0$ on A_t , which is impossible because the total charge of $\tilde{\eta}_t$ is one.

By iteration one can show a sequence $\gamma_0 > \gamma_1 > \dots > \gamma_{n-1}$ such that $g^{(m)}(x)$ is negative on $[0, \gamma_m)$ and positive on $(\gamma_m, (1+t)/2]$ for every $m = 0, 1, \dots, n-1$. This establishes our claim ($t_c = \gamma_0$).

We now can complete the proof of the theorem as follows. If $\tilde{\eta}_1$ is not a positive measure, then there is a t_1 such that the density of $\tilde{\eta}_1$ is positive on $[-1, t_1)$ and negative on $(t_1, 1]$. Then the signed equilibrium for Σ_{t_1} is given by

$$\tilde{\eta}_{t_1} = \tilde{\eta}_1^+ - \text{Bal}_s(\tilde{\eta}_1^-, \Sigma_{t_1}) - (\|\tilde{\eta}_1^-\| - \|\text{Bal}_s(\tilde{\eta}_1^-, \Sigma_{t_1})\|) \nu_{t_1} / \|\nu_{t_1}\|.$$

If it is still not a positive measure, then there exists a t_2 such that $\tilde{\eta}_{t_1}$ has positive density on $[-1, t_2)$ and negative one on $(t_2, t_1]$. Continuing the argument we derive a decreasing sequence $\{t_k\}$ with the property that $\tilde{\eta}_{t_k}$ is positive on $[-1, t_{k+1})$ and negative on $(t_{k+1}, t_k]$. The limit of this sequence is the number t_λ defined in Theorem 1.25. Thus, $t_\lambda = \max\{t : \tilde{\eta}_t \geq 0\}$, $\mu_Q = \tilde{\eta}_{t_\lambda}$, and $\text{supp}(\mu_Q) = \Sigma_{t_\lambda}$.

The Mhaskar-Saff functional \mathcal{F}_s is minimized for Σ_{t_λ} . Since $\mathcal{F}_s(\Sigma_t) = \Phi_s(t)$ (cf. Remark 1.8 and beginning of this proof), we will show similar as in the proof of

Theorem 1.14 above that t_λ is, in fact, the unique solution in $(-1, 1]$ of the relation

$$(8.3) \quad \Delta(t) := \tilde{\Phi}_s(t) - \int \frac{(R+1)^{d-s}}{(R^2 - 2Rt + 1)^{d/2}} d\lambda(R) = 0,$$

or $t_\lambda = 1$ when such a solution does not exist.

Using Quotient Rule and $\|\tilde{\varepsilon}_t\|' = d \|\tilde{\varepsilon}_t\| / dt = \int \|\varepsilon_{t,R}\|' d\lambda(R)$, we obtain

$$\tilde{\Phi}'_s(t) = - \|\nu_t\|' / \|\nu_t\| \Delta(t).$$

Observe that $\Delta(t) \rightarrow \infty$ as $t \rightarrow -1^+$. Hence, by the above relation, $\tilde{\Phi}_s(t)$ is strictly monotonically decreasing on $(-1, t')$ for some maximal $t' \in (-1, 1]$ (cf. (4.2)). If $t' = 1$, then $t_\lambda = 1$. Otherwise, $t' < 1$ and $\tilde{\Phi}_s(t') = 0$ meaning that t' is a solution of (8.3). Arguing as in the proof of Theorem 1.14 we have that every solution $t_0 \in (-1, 1)$ of (8.3) is actually a local minimum of $\tilde{\Phi}_s(t)$ because of $\tilde{\Phi}_s''(t_0) > 0$. We conclude that $\tilde{\Phi}_s(t)$ can have at most one minimum in $(-1, 1)$. Consequently $t_\lambda = t'$. We also infer that $\Delta > 0$ on $(-1, t_\lambda)$ and $\Delta < 0$ on $(t_\lambda, 1]$. This completes the proof. \square

Proof of Theorem 1.26. By definition (1.34) and Lemmas 7.2 and 7.5 one can easily see that

$$U_{d-2}^{\tilde{\varepsilon}_t}(\mathbf{z}) = \int U_{d-2}^{\text{Bal}_{d-2}(\delta_{R\mathbf{p}}, \Sigma_t)}(\mathbf{z}) d\lambda(R) = \int \frac{d\lambda(R)}{|\mathbf{z} - R\mathbf{p}|^{d-2}} = Q(\mathbf{z}), \quad \mathbf{z} \in \Sigma_t,$$

and $U_{d-2}^{\tilde{\nu}_t}(\mathbf{z}) = W_{d-2}(\mathbb{S}^d)$ on Σ_t ; hence, the weighted $(d-2)$ -potential of the signed measure $\tilde{\eta}_t$ is constant on Σ_t , that is

$$U_{d-2}^{\tilde{\eta}_t}(\mathbf{z}) + Q(\mathbf{z}) = \tilde{\Phi}_{d-2}(t) = \mathcal{F}_{d-2}(\Sigma_t), \quad \text{on } \Sigma_t.$$

The last relation follows from Remark 1.8. Moreover (with additional indication of the dependence on the parameter R),

$$\tilde{\eta}_t = \int \left[\frac{\tilde{\Phi}_{d-2}(t)}{W_{d-2}(\mathbb{S}^d)} \frac{1}{\|\lambda\|} \bar{\nu}_t - \bar{\varepsilon}_{t,R} \right] d\lambda(R).$$

Thus, the signed equilibrium is

$$(8.4) \quad d\tilde{\eta}_t(\mathbf{x}) = \left[\int \tilde{\eta}_t''(u, R) d\lambda(R) \right] d\sigma_d|_{\Sigma_t}(\mathbf{x}) + \left[\int \tilde{\eta}_t'''(u, R) d\lambda(R) \right] d\beta_t(\mathbf{x}),$$

where, when using Lemmas 7.2 and 7.5, we have for $-1 \leq u \leq t$

$$(8.5) \quad \tilde{\eta}_t''(u, R) = \frac{1}{W_{d-2}(\mathbb{S}^d)} \frac{1}{\|\lambda\|} \left[\tilde{\Phi}_{d-2}(t) - \frac{\|\lambda\| (R^2 - 1)^2}{(R^2 - 2Ru + 1)^{d/2+1}} \right],$$

$$(8.6) \quad \tilde{\eta}_t'''(u, R) = \frac{1}{\|\lambda\|} \frac{1-t}{2} \left[\tilde{\Phi}_{d-2}(t) - \frac{\|\lambda\| (R+1)^2}{(R^2 - 2Rt + 1)^{d/2}} \right] (1-t^2)^{d/2-1}.$$

It can be shown that the density with respect to $\sigma_d|_{\Sigma_t}$,

$$g(u) := \int \tilde{\eta}_t''(u, R) d\lambda(R) = \frac{1}{W_{d-2}(\mathbb{S}^d)} \left[\tilde{\Phi}_{d-2}(t) - \int \frac{(R^2 - 1)^2 d\lambda(R)}{(R^2 - 2Ru + 1)^{d/2+1}} \right],$$

is either positive for all $u \in [-1, t]$, or is positive on some interval $[-1, t_c)$ and negative on $(t_c, t]$. This follows easily from the fact that $g(u)$ is a strictly monotonically

decreasing continuous function on $[-1, t]$. Now, we turn to the density with respect to β_t , that is

$$h(u) := \int \tilde{\eta}_t'''(u, R) \, d\lambda(R) = \frac{1-t}{2} \left[\tilde{\Phi}_{d-2}(t) - \int \frac{(R+1)^2 \, d\lambda(R)}{(R^2 - 2Rt + 1)^{d/2}} \right] (1-t^2)^{d/2-1}.$$

Observe that non-negativity of the above square bracketed expression implies $g(u) \geq 0$ and therefore $\tilde{\eta}_t \geq 0$. On the other hand, if $\tilde{\eta}_t \geq 0$, then

$$(8.7) \quad \tilde{\Phi}_{d-2}(t) \geq \int \frac{(R+1)^2 \, d\lambda(R)}{(R^2 - 2Rt + 1)^{d/2}}.$$

Hence, the last relation holds if and only if $\tilde{\eta}_t \geq 0$. Note that $\tilde{\Phi}_{d-2}(t) \rightarrow \infty$ as $t \rightarrow -1^+$ and $\tilde{\Phi}_{d-2}(1) = \mathcal{F}_{d-2}(\mathbb{S}^d)$. Set $t_\lambda := \{t : \tilde{\eta}_t \geq 0\}$. Arguing as in the proof of Theorem 1.25 it can be shown that for $t = t_\lambda$ equality holds in (8.7) and t_λ is the unique minimum of $\tilde{\Phi}_{d-2}(t)$ on $(-1, 1)$ if it exists, or $t_\lambda = 1$. In particular,

$$\frac{d\tilde{\Phi}_{d-2}(t)}{dt} = -\frac{\{\|\nu_t\|\}'}{\|\nu_t\|} \left[\tilde{\Phi}_{d-2}(t) - \int \frac{(R+1)^2}{(R^2 - 2Rt + 1)^{d/2}} \, d\lambda(R) \right].$$

It remains to show the weak* convergence in (1.62). This follows from (1.36), since for any function f continuous on \mathbb{S}^d we have

$$\begin{aligned} \int_{\Sigma_t} f \, d\tilde{\epsilon}_{t,s} &= \int_{\Sigma_t} f \, d \left(\int \text{Bal}_s(\delta_{R\mathbf{p}}, \Sigma_t) \, d\lambda(R) \right) = \int_{\Sigma_t} f \, d \left(\int \epsilon_{t,s} \, d\lambda(R) \right) \\ &= \int_{\Sigma_t} f \left(\int d\epsilon_{t,s} \, d\lambda(R) \right) = \int \left(\int_{\Sigma_t} f \, d\epsilon_{t,s} \right) \, d\lambda(R) \\ &\rightarrow \int \left(\int_{\Sigma_t} f \, d\bar{\epsilon}_t \right) \, d\lambda(R) = \int_{\Sigma_t} f \, d\tilde{\epsilon}_t \quad \text{as } s \rightarrow (d-2)^+. \end{aligned}$$

This completes the proof. \square

Proof of Theorem 1.27. First observe that

$$\|\tilde{\epsilon}_{t,0}\| = \int \|\bar{\epsilon}_{t,0}\| \, d\lambda(R) = \int d\lambda(R) = \|\lambda\|,$$

which follows from principle of superposition and preservation of mass when using logarithmic balayage. Hence, $\|\tilde{\eta}_{t,0}\| = 1$ by construction. The representation (1.67) can be easily obtained using Lemmas 7.8 and 7.10. In particular, it follows from Lemma 7.10 that for $\mathbf{z} \in \Sigma_t$ there holds

$$\begin{aligned} U_0^{\tilde{\epsilon}_{t,0}}(\mathbf{z}) &= \int U_0^{\bar{\epsilon}_{t,0}}(\mathbf{z}) \, d\lambda(R) \\ &= \tilde{Q}(\mathbf{z}) + \int \left[\frac{1}{2} \log \frac{R^2 - 2Rt + 1}{2(1+t)} + \frac{(R+1)^2}{8R} \log \frac{(R+1)^2}{R^2 - 2Rt + 1} \right] \, d\lambda(R). \end{aligned}$$

By Lemma 7.8, $U_0^{\bar{\nu}_{t,0}}(\mathbf{z}) = W_0(\Sigma_t)$ on Σ_t . Hence, for $\mathbf{z} \in \Sigma_t$

$$\begin{aligned} U_0^{\tilde{\eta}_{t,0}}(\mathbf{z}) + \tilde{Q}(\mathbf{z}) &= (1 + \|\lambda\|) W_0(\Sigma_t) - U_0^{\bar{\epsilon}_{t,0}}(\mathbf{z}) + \tilde{Q}(\mathbf{z}) = W_0(\Sigma_t) \\ &+ \int \left[W_0(\Sigma_t) - \frac{1}{2} \log \frac{R^2 - 2Rt + 1}{2(1+t)} - \frac{(R+1)^2}{8R} \log \frac{(R+1)^2}{R^2 - 2Rt + 1} \right] \, d\lambda(R). \end{aligned}$$

After substituting $W_0(\Sigma_t)$ with (7.23), the above integral becomes the right-hand side of (cf. proof of Lemma 7.9)

$$\begin{aligned} \int \tilde{Q} \, d\mu_{\Sigma_t} &= \int \left(\int_{\Sigma_t} \log \frac{1}{|\mathbf{x} - \mathbf{a}|} \, d\mu_{\Sigma_t}(\mathbf{x}) \right) \, d\lambda(R) \\ &= \|\lambda\| \frac{1+t}{4} + \int \left[\frac{(R-1)^2 \log(R^2 - 2Rt + 1) - (R+1)^2 \log(R+1)^2}{8R} \right] \, d\lambda(R). \end{aligned}$$

Thus,

$$U_0^{\tilde{\eta}_{t,0}}(\mathbf{z}) + \tilde{Q}(\mathbf{z}) = W_0(\Sigma_t) + \int \tilde{Q} \, d\mu_{\Sigma_t} =: \tilde{\mathcal{F}}_0(\Sigma_t) := \tilde{\mathcal{F}}_0(t), \quad \mathbf{z} \in \Sigma_t.$$

A similar computation shows that for $\mathbf{z} \in \mathbb{S}^2 \setminus \Sigma_t$

$$U_0^{\tilde{\eta}_{t,0}}(\mathbf{z}) + \tilde{Q}(\mathbf{z}) = \tilde{\mathcal{F}}_0(\Sigma_t) + \frac{1}{2} \log \frac{1+t}{1+\xi} + \int \frac{1}{2} \log \frac{R^2 - 2Rt + 1}{R^2 - 2R\xi + 1} \, d\lambda(R).$$

The Mhaskar-Saff functional $\tilde{\mathcal{F}}_0$ for spherical caps Σ_t can be represented as

$$(8.8) \quad \begin{aligned} \tilde{\mathcal{F}}_0(t) &= (1 + \|\lambda\|) \frac{1+t}{4} - \frac{\log 2}{2} - \frac{1}{2} \log(1+t) \\ &+ \int \left[\frac{(R-1)^2 \log(R^2 - 2Rt + 1) - (R+1)^2 \log(R+1)^2}{8R} \right] \, d\lambda(R), \end{aligned}$$

which yields

$$(8.9) \quad \frac{d}{dt} \tilde{\mathcal{F}}_0(t) = \frac{1 + \|\lambda\|}{4} - \frac{1}{2(1+t)} - \frac{1}{4} \int \frac{(R-1)^2}{R^2 - 2Rt + 1} \, d\lambda(R).$$

It follows that $\tilde{\mathcal{F}}_0'(t) \rightarrow -\infty$ as $t \rightarrow -1^+$ and $\tilde{\mathcal{F}}_0'(1) = 0$. If $t < 1$, the equation $\tilde{\mathcal{F}}_0'(t) = 0$ is equivalent with each of the following two relations

$$\frac{1-t}{1+t} = \int \frac{2R(1-t)}{R^2 - 2Rt + 1} \, d\lambda(R), \quad \frac{2}{1+t} = 1 + \|\lambda\| - \int \frac{(R-1)^2}{R^2 - 2Rt + 1} \, d\lambda(R),$$

which combined give

$$(8.10) \quad 1 + \|\lambda\| = \int \frac{(R+1)^2}{R^2 - 2Rt + 1} \, d\lambda(R).$$

The above right-hand side is a strictly monotonically increasing function in t . Hence, above relation has a unique solution t_λ in $(-1, 1)$ if such a solution exists (which is a minimum of $\tilde{\mathcal{F}}_0(t)$ by the properties of $\tilde{\mathcal{F}}_0'(t)$). If there is no such solution, then $\tilde{\mathcal{F}}_0'(t) < 0$ on $(-1, 1)$ and $\tilde{\mathcal{F}}_0(t)$ is strictly monotonically decreasing on $(-1, 1)$. We conclude, that the Mhaskar-Saff functional $\tilde{\mathcal{F}}_0$ is minimized for Σ_{t_λ} , where either $t_\lambda \in (-1, 1]$ is the unique solution of equation (8.10), or $t_\lambda = 1$ if such a solution does not exist. It follows from the representation (1.67) that $\tilde{\eta}_t \geq 0$ if and only if

$$(8.11) \quad 1 + \|\lambda\| \geq \int \frac{(R+1)^2}{R^2 - 2Rt + 1} \, d\lambda(R).$$

Hence, $\tilde{\eta}_t \geq 0$ if and only if $t \leq t_\lambda$, that is $t_\lambda = \max\{t : \tilde{\eta}_t \geq 0\}$, $\text{supp}(\mu_{\tilde{Q}}) = \Sigma_{t_\lambda}$, and $\mu_{\tilde{Q}} = \tilde{\eta}_{t_\lambda,0}$. \square

APPENDIX A

Lemma A.1. *Let $-1 \leq a < b < c \leq 1$ and $|y| < 1$. Then*

$$(A.1) \quad \int_a^b (u-a)^{\beta-1} (b-u)^{\gamma-1} (c-u)^{-\alpha} {}_2\tilde{F}_1\left(\alpha, \beta; \gamma; y \frac{b-u}{c-u}\right) du \\ = \frac{\Gamma(\beta)}{\Gamma(\beta+\gamma-\alpha)\Gamma(\alpha)} (b-a)^{\beta+\gamma-1} (c-a)^{-\gamma} (c-b)^{\gamma-\alpha} (1-xy)^{-\beta} \\ \times \int_0^1 v^{\beta+\gamma-\alpha-1} (1-v)^{\alpha-1} (1-xv)^{\beta-\gamma} \left(1 - \frac{x(1-y)}{1-xy}v\right)^{-\beta} dv$$

for all $\alpha, \beta, \gamma > 0$ with $\beta + \gamma > \alpha$. Here $x := (b-a)/(c-a)$.

The last integral is the Euler type integral representation of an Appell F_1 function ([5, 5.8(5)], see also [18, 7.2.4(42)]).

Proof. A change of variable $(b-u)/(c-u) = xv$ yields

$$(A.2) \quad (b-a)^{\beta+\gamma-1} (c-a)^{-\gamma} (c-b)^{\gamma-\alpha} \\ \times \int_0^1 v^{\gamma-1} (1-v)^{\beta-1} (1-xv)^{\alpha-\beta-\gamma} {}_2\tilde{F}_1\left(\alpha, \beta; \gamma; xyv\right) dv,$$

where $0 < x < 1$. Let I denote the integral above. We substitute the series expansion of the hypergeometric function and integrate termwise.

$$(A.3) \quad I = \sum_{n=0}^{\infty} \frac{(\alpha)_n (\beta)_n}{\Gamma(n+\gamma)n!} x^n y^n \int_0^1 \frac{v^{n+\gamma-1} (1-v)^{\beta-1}}{(1-xv)^{\gamma+\beta-\alpha}} dv,$$

where the integral K_n represents the hypergeometric function ([1, Eq. 15.3.1])

$$K_n = \frac{\Gamma(\beta)\Gamma(n+\gamma)}{\Gamma(n+\beta+\gamma)} {}_2F_1\left(\gamma+\beta-\alpha, n+\gamma; n+\beta+\gamma; x\right).$$

Since $\text{Re}[n+\beta+\gamma] > \text{Re}[\gamma+\beta-\alpha] > 0$ by assumption, another application of [1, Eq. 15.3.1] gives

$$K_n = \frac{\Gamma(\beta)\Gamma(n+\gamma)}{\Gamma(\beta+\gamma-\alpha)\Gamma(n+\alpha)} \int_0^1 v^{\beta+\gamma-\alpha-1} (1-v)^{n+\alpha-1} (1-xv)^{-n-\gamma} dv.$$

Substituting the last formula into (A.3) and reversing the order of integration and summation, which is justified by uniform convergences of the series for $0 \leq v \leq 1$, yields

$$I = \frac{\Gamma(\beta)}{\Gamma(\beta+\gamma-\alpha)\Gamma(\alpha)} \int_0^1 \frac{v^{\beta+\gamma-\alpha-1} (1-v)^{\alpha-1}}{(1-xv)^\gamma} \sum_{n=0}^{\infty} \frac{(\beta)_n}{n!} \left(xy \frac{1-v}{1-xv}\right)^n dv.$$

The infinite series equals $(1-xv)^\beta (1-xv-xy+xyv)^{-\beta}$. Thus, we get

$$I = \frac{\Gamma(\beta)}{\Gamma(\beta+\gamma-\alpha)\Gamma(\alpha)} (1-xy)^{-\beta} \\ \times \int_0^1 v^{\beta+\gamma-\alpha-1} (1-v)^{\alpha-1} (1-xv)^{\beta-\gamma} \left(1 - \frac{x(1-y)}{1-xy}v\right)^{-\beta} dv.$$

We needed that $\text{Re}[\gamma+\beta-\alpha] > 0$ and $\text{Re}[\alpha] > 0$ as well as $x < 1$ and $x(1-y)/(1-xy) < 1$. The last formula for I and (A.2) give (A.1). \square

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